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Heat transfer coefficients of condensing vapors

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Earl Adolph Nirmaier

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HEAT TRANSFER COEFFICIENTS OF CONDENSING VAPORS

BY

CHARLES R. LIMA MID
EARLE A. MIRMAIRR

A THESIS SUBMITTED TO THE FACULTY OF THE DEPARTMENT OF CHEMICAL ENGINEERING œ **MEWARE COLLEGE OF ENGINEERING**

> **IN PARTIAL FULFILLMENT OF THE REQUIREMENTS FOR THE DEGREE**

> > **or**

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1956

ABSTRACT

The film coefficients of heat transfer for vapors condensing on a single horizontal tube have been measured for methyl, i-propyl, and n-butyl alcohols. The equipment used was specifically designed for these studies, and represents an improvement over that used by other investigators. The data collected were analyzed using modified Wilson method which is based on a rigorous theoretical analysis.

It has been found that the condensing film coefficients for methyl and i-propyl alcohols vary with the heat duty. This is to be expected from the analysis mentioned above. Me variation was found for n-butyl alcohol.

The observed values of the heat transfer coefficients have been compared with the values calculated with the Nusselt and Bromley equations. The Bromley equation is a modification of the Nusselt equation and takes into account the effect of the heat capacity of condensate. Excellent agreement between the observed coefficients and those predicted by the Nusselt equation were obtained with methyl and i-propyl alcohols. Good agreement was found with n-butyl alcohol. in 411 cases the Bromley equation predicted higher values of the condensing file coefficient than did the Nueselt equation.

Close agreement between the observed and theoretical coefficients is a result of the refined theoretical and experimental methods used.

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INTRODUCTION

Investigators have generally used one of two methods for measuring the file coefficients of heat transfer for vapors condensing on single horizontal tubes. These are the embedded thermocouple method as used by Kirkbride (14), Othmer and Berman (18), and Wallace and Davison (24); and the Wilson method (26). Both of these methods have inherent dis**advantages.**

Rhodes and Younger (23) pointed out that the average wall temperature as measured with thermocouples can be computed only by making certain assumptions that may not be valid. Baker and Mueller (1) **later proved that there is no point on a tube wall at which a thermocouple can be located to obtain a truly representative tube wall surface temperature.**

Rhodes and Younger (23) showed that some of Wilson's assumptions were not correct. Based en this earlier work, Chu, Plitoraft, and Solomon (6) proposed a unique modification of the Wilson method using a rigorous theoretical analysis. This modified technique has greatly improved the Wilson method.

Chu and associates' proposed technique has been applied to only a limited extent. As a result, this investigation was initiated to test this modified Wilson method using another homologous series of organic compounds. Methyl, i-propyl, and n-butyl alcohols have been selected far these studies. The physical properties of these alcohols **have been widely investigated and are more reliable than those of more complex organic compounds. Thus, the predicted** condensing **film coefficients, which are to be compared with the observed coefficients, can be calculated with a reasonable assurance of reliability.**

The predicted coefficients will be calculated using both the Nusselt and Bromley equations. The Bromley equation (4) is **a modification of the Nusselt equation and includes a correction for the effect of the heat capacity of condensate. The Bromley equation is to be compared with Nusselt's equation to determine the magnitude of this correction.**

The equipment used to measure the beat transfer coefficients vas specifically designed for these studies and represents an Improvement over that used by the other investigators.

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THEORY

The Nusselt equation (15) (17) is generally used to predict the coefficients of heat transfer fat the filmdom condensation of pure vapors on a colder surface. Applied to a single, horizontal, cylindrical tube, the equation is

$$
h_0 = 0.725 (k_f^3 / \rho_f^2 g) / D_0 u_f \Delta t)^{0.25}
$$

which is based en the assumption that streamline flow exists throughout a continuous condensate film, and that gravity alone causes the condensate to flew over the smooth surface. The possible acceleration effect of the vapor velocity upon the film thickness is also neglected.

Many investigators (14) (16) (24) have measured the condensing film coefficients of pure organic) vapors with the aid of embedded thermocouples. inasmuch as the condensate film thickness around the perifery of a cylindrical tube varies, it is to be expected that the file temperature would also very. Therefore, the measurement of surface temperatures with thermocouples may result in considerable error. Baker end Mueller (1) proved that there is no point on the surface of a cylindrical tube at which a thermocouple can be located to Obtain a representative surface temperature.

Instigators (3) (23) have utilised the Wilson method to avoid the difficulties associated with direct temperature measurements. The Wilson method is represented by the following equation

$$
\mathbf{R} = \sum_{i=1}^{n} \mathbf{R}_{i} + \mathbf{R}_{i} + a/\nabla^{0.8}
$$

where Rv is the thermal resistance of the condensate film at a finite cooling water flow rate, and "0 is a constant. This equation is based on the assumption that changes in cooling water rate have no effect upon R_u. Rhodes and Younger (23) pointed out that R_u varied **with water rate and postulated that**

$$
R_{\mathbf{v}} = R_{\mathbf{v}_0} + b/\mathbf{v}^{0.8}
$$
 III

Beatty and Kats (3) applied this method to their work with finned tubes, with a modification to allow for the effect of cooling water temperature upon the water film resistance.

Chu, Flitcraft, and Reisman (6) **stated that Equation III was empirical,** *since* **there is no theoretical justification far setting up R_v** as a function of R_{vo} and the cooling water *flow* rate. They **proposed a modification of the Wilson method based on a rigorous theoretical analysis. It was pointed out in their paper and by other 1nvestigators (7) (23) that the group** $(k_f^3 \int_f^2 g_A / \mu_f)^{0.25}$ **appears to remain constant with** *varying* **temperature for most organic solvents. Therefore, for steady state heat transfer through a condensing vapor, the following equation has been derived** (6).

$$
h_{\phi} = c q^{-1/3}
$$

This equation will result in a straight line with a negative slope equal to one-third on a log—log plot. Thus, Chu and associates (6) **have deduced that the condensing film coefficient varies with the heat transferred.**

The values of the group $(k_f^3 f_f^2 g_A/M_f)^{0.25}$ for methyl, i-propyl,

and n-butyl alcohols are listed in Table 1. Also included in this listing are the values of this group for other organic compounds as reported by Chu and associates (7).

The overall thermal resistance from the condensing vapor to the cooling water side of a condenser tube is represented by

D;3•2 1 = 1 x VI; Ccr+1135-7Pi717- -2"MTI°.8

From Equation IV it is evident that h_a is constant at constant values **of q. Also, the thermal resistance of the tube wall is negligible compared to the other terms of Equation V. Therefore, the first two terms of Equation V are constant at constant values of q. A plot of** $1/\mathbb{U}_{\alpha}A_{\alpha}$ against $1/(1 \ 0.011t)\sqrt{0.8}$ should result in a straight line at equal values of q. The intercept of this line, equal to l/h_ok_o x/kwA_{av}, **can be used to calculate the condensing film coefficient, he, since**

$$
\frac{1}{\sigma_{\alpha} \sigma_{\alpha}} = -\frac{1}{\alpha} = \frac{1}{\sigma_{\alpha} \sigma_{\alpha}} + \frac{1}{\sigma_{\alpha} \sigma_{\alpha}}
$$

and

$$
\frac{1}{h_0} = \frac{1}{h_0} - \frac{1}{h_0} \sum_{\text{av}} f_0
$$

Several sets of experimental data can be obtained by varying the cooling water flow rate and the pressure in the test condenser vapor space. For any one set, the overall water to vapor temperature difference would be kept constant; and from one set to another the temperature difference changed by readjusting the pressure in the vapor space. Two or more points, at constant q, can be gotten from the above sets of data. These points, when replotted as $1/U_{ab}A_{0}$ against

TABLE 1

RATIO OF $(\mathbf{k_f}^3 f_g^2 \mathbf{s_M}/\mu_f)^{0.25}$ at different temperatures

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 $1/(1+0.011t) v^{0.8}$, yield a straight line. The value of h_0 at the **constant value of q is then calculated from the intercept of this** $11ms_•$

As previously stated, the Musselt equation is used to predict **condensing film coefficients. Bromley (4) has pointed out that in** the basic derivation by Musselt the effect of the heat capacity of **condensate was assumed negligible. This assmuption is valid at law pressures. Weever, at high pressures and large tube wall to saturated** vapor temperature differences, sensible heat may significantly affect **hest transfer.**

This consideration led to the theoretical development of a modified Nusselt equatiob. The equation for vapors condensing on a single horizontal tube is 2 11.25 \triangle **te_m** $)$ ² **1.728** $\frac{1}{2} \int_{e}^{2} \frac{\int_{e}^{2} \left(\int_{e}^{e} f(x) \right) \left(\frac{1}{2} + 0. \frac{1}{2} \right)^{2}}{\int_{e}^{2} \left(\int_{e}^{e} f(x) \right)^{2}}$ **VIII IleMf**

Bromley's equation should predict higher values of h_e than the **Nusselt equation.**

DESCRIPTION OF APPARATUS

The equipment used in thin work, se sham in Figures 1 and 2, consisted basically of a kettle, a single tube horizontal condenser, **a cooling water circulation system, and a vacuum pump.**

The jacketed kettle was of stainless steel construction and had a five gallon capacity. Steams at an initial pressure of 90 psig., was introduced into the jacket through a 3 to 15 psig. throttling valve.

The test condenser was a 0.375 inch outside diameter brass tube 24 inches long having a wall thickness of 0.035 inches. The effective outside surface area was 0.196 sq. ft., and the thermal conductivity vas 6o **Btu/(Hr.) (sq** rt.) (°F/ft.). A 2.5 **inch schedule 40 stainless steel pipes, flanged at both ends, was used as the condenser jacket.**

Taper from the kettle entered the jacket through three 0.5 inch diameter distributors. Condensate returned to the kettle by gravity flow through two downcomers and liquid seal traps. The temperature **et the uncondensed vapor taken overhead was measured by a -1 to 101°C or 99 to 201°C thermometer graduated in 0.1°C increments. A glass secondary condenser was used to assure that excess vapor was passing the test condenser at all times. Condensate returned to the kettle through a liquid seal trap.**

The condenser jacket contained a special thermocouple installation used tar measuring bulk condensate temperatures. A finely threaded fitting permitted adjustment of the position of the thermocouple junction.

CONDENSER JACKET EXPERIMENTAL APPARATUS

JO

The system was evacuated and maintained at law pressures by a Conco-Hypervac it pump of 1.44 cubic ft. per minute free air displacement. The pressure was controlled at the desired level by a cartesian manostat and measured with a U-tube mercury manometer.

The kettle - condenser - vacuum system was checked for leaks. It was considered to be air-tight when no noticeable increase in absolute pressure occurred over a one-half hour period, after evacuation to an absolute pressure level of one inch of mercury.

Water was circulated by a centrifugal pump rated at 25 gallons per minute at 80 ft. head and 1.0 specific gravity. Flow rates were measured with a calibrated Fisher-Porter rotameter rated a 13.6 gallons per minute.

Two 55 **gallon drums were used as holdup and constant head supply for the pump. Fresh water flowed to the pump from the supply drum. Heated water from the condenser could be discharged to the second drum and then to the sewer. The secondary drum was equipped with s steam coil which allowed pre-heating of the cooling water.**

Condenser inlet water temperatures were measured with a 0 to 50°C thermometer graduated in 0.1°C increments. The outlet water temperature measurements were made with a thermocouple installation. A pair of 0 to 6°C Beckmann thermometers, each of which could be read to 0.002°C, was used to determine the water temperature rise in the test condenser.

The entire kettle-condenser system was well insulated. Magnesia

lagging was used on all pipe lines. The condenser Jacket was enclosed in a plywood box which was filled with powdered vermiculite. This method of insulation minimised heat losses.

Eighteen gage, copper-constantan thermocouples were used. These thermocouples were factory calibrated at ±0.8°C up to a temperature et 375 °C.

The alcohols used were reagent grade containing a maximum of **0.5 per cent impurities. The boiling ranges were less than 1°C.**

EXPERIMENTAL PROCEDURE

The test alcohol was charged to the kettle after the system had been cleaned. Before heating with steam, the cooling water was turned on and the system was allowed to reach a steady state condition. At this point a series of reference readings was taken on the Beckmann thermometers. The rotameter reading and water temperature were also taken. This blank correction on the Beckmann thermometers was used to determine the true water temperature rise in the condenser tube, as discussed in the appendix section *Reference Correction of Beckmann Thermometers".

The system was then evacuated, bypassing the manostat to permit more rapid evacuation. When the absolute pressure decreased to the vapor pressure of the alcohol in the kettle, the pump was turned off and the bypass closed. Stem was then admitted to the kettle jacket. When the system pressure reached the desired level, the manostat was put in operation and the vacuum pump restarted. This method of startup minimised the passibility of non-condensable gases remaining in the system.

Five complete sets of readings were taken at one minute intervals after the system reached equilibrium. The average of these readings **was used in the calculations. A set of data consisted of a rotameter reading, absolute inlet and outlet water temperatures, vapor temperature, bulk condensate temperature, kettle pressure, and a simultaneous reading of the Beckmann thermometers. The secondary**

condenser wee constantly checked to assure that excess vapors were present in the test condenser at all times.

After the completion of a run, the water rate was changed. If **necessary, the pressure level was readjusted to maintain a constant overall water to vapor temperature difference.**

Twenty to thirty minute time intervals were required for the system to return to equilibrium after conditions were changed.

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EXPERIMENTAL RESULTS AND TREATMENT OF DATA

The experimental results for methyl, i-propyl, and n-butyl alcohols are presented in Tables 2, 3, and 4, which include the calculated results of (1×10^3) **/** $(1+0.011t)\mathbb{V}^{0.8}$ **.** The data are plotted in Figures 3 to 11. Figures 3, 6, and 9 show the variation **of the water flew rate and temperature rise in the test condenser. Figures 4, 7, and 10 are a form of the Wilson plot and represent** the heat transferred as a function of $1/\sqrt{9.8}$ at varying overall **vapor to cooling water temperature differences.**

In determining the heat transfer coefficients, ho, a series of **straight lines were drawn parallel to the abscissae of Figures 4, 7, and 10, intersecting two or more of the constant vapor - water temperature difference lines. These horizontal lines represent constant** values of the heat transferred, q_a . At each intersection, the overall **temperature difference, and** $(1 \times 10^3) / (1 + 0.011t)\sqrt{0.8}$ **were read.** $\frac{\Delta T}{q}$ was calculated and plotted against $(1 \times 10^3) / (1 + 0.011t)\n^0$ ^{0.5} **for each value of 4. These plots, shown in Figures 5, 8, and 11, yielded straight lines which were extrapolated to the ordinates.** The intercepts represent ∆**T/q or l/U_oA_o** at infinite water flow. The **values of he, the condensing film coefficients, are calculated from them intercepts.**

TABLE 2
TABULATED RESULTS - METHYL ALCOHOL

 $\hat{\boldsymbol{\delta}}$

TABLE 2 (Con't) TABULATED RESULTS - METHYL ALCOHOL

RUN NO.	INLET WATER TEMP. $\delta_{\rm C}$	WATER TEMP. R_{σ}^{JSE}	VAPOR TEMP $^{\circ}$ C	OVERALL WATER- VAPOR ΔT , $^{\circ}C$	WATER FLOW W LB/HR	HEAT DUTY Q BTU/HR	1x103y -. 8 1/0.011t
333890	6.50 6.55 6.55 6.60 6.75	1.307 1.440 1.633 1.696 2.137	60.76 60.76 60.76 60.75 60.80	53.61 53.49 53.39 $\frac{53.30}{52.98}$	3330 2970 2590 2345 1950	7800 7700 7630 7165 7490	1.005 1.118 1.243 1.342 1.550
41 1444万	6,90 7.20 7.40 7.60 7.80	2.478 2.983 $3 - 357$ 3.897 4.330	61.87 61.89 62.09 62.56 62.69	$53 - 73$ $53 - 20$ 53.11 53.01 52.72	1660 1300 1150 970 840	7410 6985 6950 6805 6545	1.757 2.11 2.32 2.62 2.92

 $\hat{\mathbf{v}}$

TABLE 3 **TABULATED RESULTS - ISOPROPYL ALCOHOL**

 $\mathcal{L}^{\mathcal{L}}$

 $\mathcal{L}(\mathbf{x}) = \mathcal{L}(\mathbf{x}, \mathbf{y}) = \mathcal{L}(\mathbf{x})$ $\sigma\tau$

TABLE 4 D RESULTS **- N-BUTYL ALCOHOL**

RUN NO.	INLET WATER TEMP.	WATER TEMP. RISE C	VAPOR TEMP. $\mathbf{O}_{\mathbf{C}}$	OVERALL WATER- VAPQR ΔT , $^{\circ}C$	WATER FLOW W LB/HR	HEAT DUTY Q <u>BTU/HR</u>	$1x10^3r^{-8}$ 1f0.011t
84 85 86 $\begin{array}{c} 87 \\ 88 \end{array}$	7.50 8,10 7.30 7.20 7.02	2.251 3.378 1.958 1.685 1.425	95.14 95.86 95.45 94.01 94.01	86.51 86.07 87.17 85.97 86.28	1400 930 1680 1980 2345	5665 5660 5920 6010 6020	2.00 2.72 1.74 1.53 1.34
89 90 91 92 93	7.00 7,18 9.00 8.90 7.72	1.255 1.030% 4.399 4.358 2.484	94.06 93.45 97.25 97.14 99.27	86.43 85.75 86.05 86.06 90.31	2680 3235 700 720 1365	6050 6000 5540 5650 6100	1.20 1.04 3.34 3.26 2.03
94 95 96 97 98	8.25 8.88 7.50 7.58 8.04	3.570 4.427 $1.057 -$ 1.331 2.250	99.53 $100 - 23$ 98.65 98.55 98.61	88.69 89.14 90.62 $90 - 30$ 89.44	955 750 3240 2595 1515	6140 5980 6170 6220 6140	2.62 3.16 1.03 1.23 1.86
99 100 101 102 103	7.80 7.96 8.11 8.37 8.80	1.688 1.623 1.866 2.304 2.635	99.10 104.07 104.10 104.40 105.78	90.46 95.30 95.06 94.88 95.66	2040 2265 1940 1595 1400	6210 6630 6520 6610 6630	1.47 1.36 1.53 1.78 1.96
104 105 106 107 108	9.30 9.50 6.48 8.00 8,00	3.782 4.253 5.371 1.169 1.299	106.60 107.07 103.85 103.65 103.80	95.41 95.44 94.68 95.07 95.15	960 830 650 3210 2840	6540 6350 6290 6750 6640	2.59 2.89 3.64 1.03 1.13
109 110 111 112 113	8,00 7.98 $8 - 30$ 8.35 8,50	1,420 1.220 1.714 1,901 2.422	108.67 108,38 109.35 109.79 110.45	99.96 100.40 100.19 100.49 100.74	2740 3210 2270 1980 1620	7010 7050 7000 6790 7060	1.17 1.03 1.35 1.50 1.75
114 115 116 117 118	8.80 6.67 8.23 8.30 8.80	2.992 5.517 1.237 4.605 3.406	110.44 108.60 108.78 111.31 111.07	100.49 99.17 $99 - 93$ 99.71 100.57	1310 665 3210 800 1120	7050 6610 7150 6640 6860	2.06 3.63 1.03 3.03 2.31

1 AD LE ÷ LOU LI TABULATED RESULTS - N-BUTYL ALCOHOL									
RUN NO.	INLET WATER TEMP. \circ_{C}	WATER TEMP. RISE $^{\circ}$ c	VAPOR TEMP. $\mathbf{0}_{\mathbf{C}}$	OVERALL WATER- VAPOR ΔT , $^{\circ}C$	WATER FLOW ₩ LB/HR	HEAT DUTY Q BTU/HR	1X103v-08 1/0.011t		
119 120 121 122 123	8.80 8.40 8.60 8.80 $9 - 30$	3.658 2.137 2.659 3.180 4.496	116.25 116.01 116.14 116.71 117.37	105.62 106.54 106.21 106.29 105.82	1115 1995 1580 1310 990	7340 7670 7560 7500 7370	2.33 1.49 1.78 2.05 2.69		
124 125 126	7.86 $7 - 43$ 8.81	5.768 5.479 5.056	116.79 116.27 116.54	106.00 106.10 105.20	665 705 775	6910 6950 7060	3.53 3.37 3.06		

TABLE 4 (Con't)

 $\sim 10^{-1}$

Figure 11 x is the tween 358.14 \cdot

DISCUSSION

It has been found that the value of h o varies with q for methyl and i-propyl alcohols. No variation was found for n-butyl alcohol. The values of h o far methyl alcohol ranged from 485 to 527 with q varying from 7750 to 5750. For i-propyl alcohol, he varied from 221 to 250 with q changing from 4700 to 3800. The value of h_a for **n-butyl alcohol was 208 with q ranging from 7150 to 5960.**

The h_e values were determined from lines drawn through two or **three points. Most of the values for methyl alcohol were obtained from three points and most of those for the other alcohols from two points. It is desirable to obtain as many points as possible. Rowever, this is not always practicable. As is evident from Figures 4, 7, and 10, the number of points determined is dependent upon the slopes and spacing of the curves. The slopes of the curves are a function of each individual alcohol; the spacing is an experimental variable.**

The data of Chu, et. al.; (7) indicate that the slopes of the above mentioned curves become very steep at high cooling water flew rates. Pump head limitations in this work prevented the experimental determination of the steep end of these curves. Investigation at high water rates would have provided additional points far the determination of h_a.

Reducing the spacing between curves to obtain more points was

inpracticable due to data overlap. This overlap is best demonstrated Figures 3, 6, and 9.

it would have been desirable to make investigations over wider ranges of q. Limitations of the pressure controller made it impossible to obtain lower values of this parameter.

Since many of the values of h_o were determined from only two **points, it is necessary to discuss the accuracy of these pairs of points. For the moat part, these points were obtained from the head of one curve and the tail of another. (Figures 4, 7, and 10) The head, or high cooling water rate end, was usually well defined by a large number of data points. The tail, or low water rate end, was usually defined by fewer points, tending to make this end of the curve less reliable. Also, as the overall temperature difference decreased, the curves became flatter. This wee to be expected since from a theoretical consideration, q will equal sore when the overall temperature difference is sere, no matter what the rate of cooling water flow. A slight change in elope of these flat curves can make** a large difference in the intercept selected to calculate h_{ab} . **Therefore, the points obtained from low water rates and/or low overall temperature differences are least reliable.**

The water flow rate variation was less than one per cent, except at low flows where it was occasionally as high as three per cent, It is believed, therefore, that the value of q is correct to within ±3 per cent. The Beckmann thermometers could be read to within

0.002°C. Since the minimum cooling water temperature rise was about 0.800°C the maximum error involved in reading these thermometers is ±0.5 per cent. The average variation of the overall water to vapor temperature difference is ±1 per cent. This variation was kept to a minimum by adjusting vapor temperature as water temperature changed. The overall error of the observed condensing film coefficients is estimated to be 5 to 10 per cent.

As shown in a previous section, the relation of h_e and q is such that a log-log plot of h_a and q should give a straight line with a **slope of minus one-third. This plot is shown in Figure 12 and the data listed in Table 5. The data on methyl and i-propyl alcohols confirm this theoretical relationship. The slope for methyl alcohol was -0.278 and for i-propyl alcohol was -0.405. The deviation of these slopes from -0.333 is within the accuracy of the graphical techniques involved. Me variation of he and q was found for n-butyl alcohol. The reasons for this are unknown at this time.**

The value of h_o at the lowest measured value of **q** for i-propyl **alcohol is inconsistent with the other determinations. As discussed above, this inconsistency is probably a result of the inaccuracies** inherent in the determination of h_o at low values of q.

As discussed in the "Theory", Bromley re-derived the Nusselt equation, taking into account the effect of the heat capacity of condensate. The observed results of the condensing film coefficients of heat transfer have been compared to those predicted by both the

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Nusselt and Bromley equations. These results are presented in Table 5.

The agreement between the observed values of h_o and those predicted **by the Nusselt equation is excellent for methyl and i-propyl alcohols.** The ratio of h_{obs}/h_{calc} ranges from 0.975 to 1.020 for methyl alcohol **and from 0.975 to 0.992 far i-propyl alcohol, with a ratio of 0.863 for the one inconsistent data point for i-propyl alcohol. The agreement between the observed and the predicted values of n-butyl alcohol is good, the ratio ranging from 0.879 to 0.934. For the most part,** for the alcohols studied, the observed values of h_e were lower than **the theoretical values.**

The condensing film coefficients predicted by the Bromley equation are higher than those predicted by the Nusselt equation. Bromley's predictions were 2% higher for methyl alcohol, 4 to 5 per cent higher for i-propyl alcohol, and 6 to 7 per cent higher for n-butyl alcohol. This equation, although supposedly better than Nusselt's at high pressures, apparently overpredicts coefficients at low and reduced pressures. It can be concluded from these observations that the Nusselt equation satisfactorily predicts condensing film coefficients at low pressures where the effect of the heat capacity is negligible.

The values of h_e were calculated at constant values of q. Since **the variations of latent heats were small over the range investigated in this work, the tube loadings and film thicknesses for any one** alcohol remained essentially constant for these values of h_o. Also, **the acceleration effects of the vapor on the condensate film were**

minimised by using a large condenser jacket. These are further reasons for the agreement between the observed and theoretical results predicted by the Nusselt equation.

This study shows that the ratio of h_{obs}/h_{calce} decreases as the **molecular weight of the alcohol increases. The average of the ratios for the three alcohols in order of increasing molecular weight are 0.997, 0.950, and 0.910. It should be noted that the work of Chu,** et. al., (7) indicates a ratio of 0.420 for n-hexyl alcohol.

It has been suggested in the past that sub-cooling of condensate occurs and effects the experimental results. In order to measure the effect of sub-cooling, if any, a thermocouple was installed to measure the condensate temperature. These temperature data are listed in Tables 8, 9, and 10.

In all cases the measured condensate temperatures were about 5 to 10°C lower than the saturated vapor temperatures. Since this measured bulk condensate temperature is higher than any average calculated with the equation recommended by Mc Adams (15) , $t_f = t_{gy} -0.75$ $(t_{gy} - t_g)$, **it is felt that sub-cooling did not effect the experimental results of this work. Furthermore, excess vapors were present in the test condenser at all times.**

Another possible source of error is the presence of noncondensable gases in **the condensing vapor. The experimental technique, as described in a previous section, precludes this possibility.**

The fouling of a condenser tube surface effects the condensing film coefficients. Rhodes and Younger (23) found that the apparent **thermal resistance of a fouled tube can be 10 to 20 per cent higher than that of a cleaned tube. This may partially account for the fact** that the ratios of h_{obs}/h_{oale} were usually less than unity.

The predicted condensing film coefficients are always affected by the physical property data used. The properties of the alcohols used in this work were obtained from the International Critical Tables (10) and from other, mare recent, literature (2) (8) (9) (15) (21). There is considerable inconsistency in these data, especially far thermal conductivity.

 ~ 100 km s $^{-1}$

FIGURE 12

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CONCLUSIONS

The technique of Chu, Flitcraft, and Holman, for measuring tho film coefficients of condensing vapors, has been checked with methyl, i-propyl, and n-butyl alcohols. Chu and associates concluded that

$$
h_{\bullet} \equiv e \cdot q^{-1/3}
$$

Tho work on methyl and i-propyl alcohols presented in this paper substantiate their conclusion. No variation of h_e and q was found **for n-butyl alcohol.**

The theoretical condensing film coefficients calculated from the Nusselt and Bromley equations were compared with the observed values. The Musselt equation predicted values of h_e closer to those **Observed than did the Bromley equation. Excellent agreement of theoretical and observed coefficients was found for methyl and i-propyl** alcohols, and good agreement was found for n-butyl alcohol. In all cases the Bromley equation predicted higher values of h_e than did the **Nusselt equation. It can be concluded from these observations that the Nusselt equation satisfactorily predicts condensing film coefficients at low pressures where the effect of the heat capacity of condensate is negligible.**

The close agreement between the theoretical and observed coefficients is attributed to 1) the rigorous theoretical analysis of the Wilson method, as presented by Chu and associates, eliminating variables not previously considered and 2) the use of equipment, designed specifically for this work, which represents an improvement over apparatus used by **other investigators.**

RECOMMENDATIONS

The following recommendations are forwarded:

- **1. Continue investigations with the aliphatic alcohols. Determine the relationship, if any, between the ratio** of h_{obs}/h_{oale} and molecular weight.
- **2. Investigate other homologous series of organic compounds.**
- **3. Conduct investigations over a wider range of cooling water flow rates. Initiate studies at positive pressures.**

NOMENCLATURE

 $\Delta \sim 10^{11}$ m $^{-1}$.

 $\mathcal{L}^{\text{max}}_{\text{max}}$

APPENDIX

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APPENDIX CONTENTS

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REFERENCE CORRECTION OF BECKMANN THERMOMETERS

Two Beckmann thermometers were used to measure the cooling water temperature rise. No attempt was made to adjust the mercury camas to the same scale reading. However, both mercury levels were adjusted to the lower ends of the scales at the prevailing cooling water temperatures for maximum utilisation of the scales.

Since the mercury columns in the Beckmann thermometers were not preset to the same scale reading, the indicated temperature difference at sere heat duty had to be determined. This was accomplished by pumping water through the condenser tube at a constant rate until the system reached equilibrium. The temperature difference was then read. This reading, or "blank", was algebraically added to the **indicated temperature rise at a finite heat duty to obtain the true increase.**

It was found that there was a small, but measurable, rise in water temperature due to friction in the condenser tube. This temperature rise increased with increasing water flow rate. A calibration curve was developed to determine the magnitude of this effect. This curve is shown in Figure 14. It is evident that this temperature rise would have a significant effect upon temperature differences measured under test conditions. Therefore, a correction for the friction effect, as shown in Figure 14, was employed.

It is convenient to illustrate the use of the correction factors

with an sample. Assume that a "blank" was established for the Beckmann thermometers at a rotameter reading of 50 per cent. At this water rate, a temperature difference of 0.500°C wee indicated. Under test conditions at the same flow rate, say the temperature difference was 1.700°C. The actual temperature difference in the condenser tubs would then be 1.700 - 0.500 or 1.200°C. As shown by Figure 14, the increase due to friction alone was 0,262°C. Therefore, the true temperature difference resulting from the condensing vapor is 1.200 - 0.262°C or 0.938°C.

All differential temperature data presented in this work were corrected in the above manner.

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TABLE 8 **ORIGINAL DATE - METHYL ALCOHOL**

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RUN NO.	ROTA- METER RDG	INLET WATER TEMPERATURE		OUTLET WATER TEMPERATURE		VAPOR TEMP.		FILM TEMP.	
	≰	$ABSO-$ LUTE $\circ_{\textsf{C}}$	BECK- MANN $\mathbf{P}_{\mathbf{C}}$	$ABSO -$ LUTE NY.	BECK- MANN \mathbf{C}	$\circ_{\mathbb{C}}$	MV	$\mathbf{^o}_\mathbf{C}$	
$\frac{25}{26}$ $\mathbf C$	19.0 14.1 50.0	7.12 7.65 $6 - 30$	3.108 3.657 2.415	0.346 0.400	4.140 5.536 0.633	52.73 52.96	1.98 2.01	48.8 49.5	
$\frac{27}{28}$	50.6 44.2	6.40 6.37	2.422 2.375	0.240 0.246	1.664 1.802	45.71 45.65	1.62 1.63	40.2 40.4	
29 30	39.8 35.7 30.9	6.40 6,50	2.365 2.430	0.248 0.258	1.936 2.157	45.79 45.75	1.65 1.65	40.9 40.9	
$\mathbf C$ $\frac{31}{32}$	$\frac{30.9}{24.8}$	6,80 $6 - 80$ 7.02	2.940 2.750 2.973	0.280 0.300	1.361 2.714 3.328	46.36 46.63	1.72 1.74	42.6 43.1	
$33 + 35$ C	$20*0$ 14.2 $50 - 3$ 50.0	7.43 8,20 6,35 6,50	$3 - 336$ 4.133 2.360 2.352	$0 - 333$ 0.400 0.250	4.088 5.663 1.640 0.692	47.28 48.03 46.23	1.77 1.81 1.71	43.8 44.8 42.4	
36	$50 - 3$	6.50	2.417	0.256	2.061	60.76	2.27	55.4	
$\frac{37}{38}$ $\frac{39}{40}$ $1 + 1$	45.0 39.3 35.6 29.6 25.2	6.55 6.55 6.60 6.75 6.90	2.432 2.408 2.468 2.579 2.709	0.264 0.270 0.276 0.300 0.320	2.256 2.473 2.624 3.217 3.714	60.76 60.76 60.75 60.80 61.87	2.27 2.27 2.27 2,28 2.32	55.4 $\frac{55.4}{55.1}$ 55.7 56.6	
42 $\frac{c}{c}$	19.8 $\frac{13}{31}$.	$7 - 20$ $7 - 70$	3.038 3.468	0.352	4.577 1.621 0.465	61.89	$2 - 33$	$56 - 8$	
$^{43}_{44}$ 45	17.5 14.8 12.9	7.60 7.40 7.60 7.80	$2 - 399$ 3.172 $\frac{3.367}{3.540}$	0.394 0.412 0.440	4.667 5.413 6.025	62.09 62.56 62.69	2.36 2.39 2.40	$52 - 5$ $58 - 2$ 58.4	

TABLE 8 **(con't) ORIGINAL DATE - METHYL ALCOHOL**

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TABLE 9 -------- \cdots

RUN NO.	ROTA- METER RDG ₫	INLET WATER TEMPERATURE		OUTLET WATER TEMPERATURE		VAPOR TEMP.		FILM TEMP.
		ABSO- LUTE $^{\circ}$ C	$BECK-$ MANN $^{\circ}$ c	ABSO- LUTE MY	$BECK-$ MANN $^{\circ}$ C	\mathbf{c}	NY	$^{\circ}$ c $-$
72 73	42.2 48.7	6.10 6,00	1.748 1.714	0.230 0,220	0.705 0.509	64.41 64.74	2.40 2.40	58.4 58.4
\mathbf{C} $\frac{74}{75}$	41.0 $35 - 0$ 29.3	5.90 5.90 6.00	1,468 1.475 1.580	0.230 0.240	1.405 2.372 2,697	59.09 59.46	2.20 2,20	53.9 53.9
76 $\frac{77}{78}$ 79 $\mathbf C$	41.0 49.0 23.4 19.3 18.9	5.80 5.75 6.24 6.50 6,70	1,366 1.363 1.752 1.994 2.251	0.210 0.210 0.250 0.275	2.077 1.867 3.138 3,653 0.495	58.28 57.76 59.98 60.44	2.16 2.15 2.24 2.26	$53 - 0$ $\frac{52.7}{54.8}$ 55.2
80 81 82 83	16.3 12.7 7.8 10.2	6,87 7.20 8.40 7.63	$2 - 464$ 2,823 4.098 3.225	0.300 0.340 0.445 0.375	2.665 3.531 6.117 4.384	61.09 61.51 62.44 61.58	$2*31$ 2.32 2.36 2.33	56.3 56.6 57.5 56.8

TABLE 9 (Con't)
ORIGINAL DATE - ISOPROPYL ALCOHOL

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TABLE 10 (con't)

SAMPLE CALCULATIONS

I. RUN NO. 2 - METHYL ALCOHOL

1. Rate of Heat Transfer, q

Water flow rate \pm 1950 lb./hr. Water temperature rise = 2.280° C Heat depacity of water = 1.0 Btu/(1b) (°F) $q = 1950$ (1.0) (2.280) (1.8) = 8010 Btu/hr.

Inlet water temperature $p = 7.60^{\circ}$ C Water temperature rise = 2.280° C

Bulk temp. ($^{\circ}$ C) = 5.80 2.280 = 6.94 $^{\circ}$ C

Bulk temp. $({}^{\circ}F) = (6.9h) (1.8) + 32 = h h.6^{\circ}F$

3. Overall Water to Vapor Temperature Difference, °C Water bulk temperature = $6.9h^{\circ}$ C Saturated vapor temperature \pm 64.60°C

Overall temp. difference \pm 64.60 - 6.94 \pm 57.66 °C

$$
h. \frac{1 \times 10^3}{(1+0.011 \epsilon)\sqrt{10}}.
$$

Water bulk temperature = $h\mu_0$.6 \hat{r}

Water flow rate, $V = 1950$ lb./hr.

$$
\sqrt{9} \cdot \sqrt{8} = h25
$$

$$
\frac{1\times10^{3}}{(1+0.011\epsilon)}\sqrt{0.8} = \frac{1\times10^{3}}{[1+0.011(l\mu,6)] (425)} = 1.58
$$

II. METHIL ALCOROL AT q = 7h00 Btu/hr

1. Cheerved Heat Transfer Coefficient, h.

From Figure 4 at $q = 7400$ Btu/hr

 $\frac{1 \times 10^3}{(1+0.011\epsilon)^{9.8}}$ is plotted in Figure 5.

at $\frac{1 \times 10^3}{(1+0.011t)^9}$ 0.8 = 0, 0.01025

Tube wall thickness, $x = 0.00292$ ft. Outside tube surface area = 0.1962 sq.ft. Inside tube surface area \pm 0.1596 sq.ft. Tube thermal conductivity, $k = 60$ Btu/(hr) (sq.ft.) (*F/ft)

Average tube surface area, $A_{\text{av}} = \frac{0.1962 - 0.1596}{10.1262} = 0.1772$

 $\frac{1}{\Gamma}$ = $\frac{0.00292}{60(0.1772)}$ = 0.00027 $\frac{1}{1000} = \frac{\Delta T}{9} = \frac{1}{1000} + \frac{X}{K}$ $h_0 = \frac{1}{0.19570.01025 - 0.00027} = 511 Btu/(hr) (sq_*ft_*)(^0r)$

Theoretical Heat Transfer Coefficient, h. $2.$

At infinite water flow rate, water bulk temp = inlet water temp = 7.0°C (Pigure h at $T = 53.25$ °C) Saturated vaper temp., $t_{\text{avg}} = 61.5^{\circ}C$ Thermal resistance from water to vapor ≥ 0.01025 Thermal resistance of tube wall \pm 0.00027 Water to saturated vapor temp. diff. \pm 61.5-7.0 \pm 54.5°C Tube surface temp., $t_{\text{m}} = 7.0 + 54.5$ $\left[\frac{0100027}{0.01025}\right]$ = 8.h% Film temp., $t_{f} = t_{av} - 0.75$ $(t_{av} - t_{u})$ $= 61.5 -0.75 (61.5 - 8.4) - 21.7$ [®]C Temp. drop across condensate film \pm 61.5 -8.h \pm 53.1°C At $t_r = 21.7$ ⁰C Thermal conductivity, $k_{\text{f}} = 0.1235$ Btu/(hr)(sq.ft.)($^{9}F/ft$) Liquid density, $\int_{e} = h9.2$ lb/cu.ft. Latent heat, $\lambda = 503$ Btu/1b. Viscosity, $\mu = 1$, ho lb/(hr)(ft.) Hest capacity, ep = 0.601 Btu/(1b) (σ_F) Gravitational constant, $g = h \cdot 17 \times 10^8 \text{ ft/hr}^2$ Outside tube diameter, $D_{\text{m}} = 0.03125$ ft. Temp. drop across condensate $f11m = 95.7^{\circ}F$

(a) h_a by **Musselt** equation $h_0 = 0.725 \sqrt{\frac{k_f^3 \beta_f^2}{D_{\mu \mu} \Delta t}}$

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$$
h_0 = 0.725 \sqrt{\frac{(0.1235)^3 (h_9.2)^2 (h.17 \times 10^8)}{0.03125 (1.10) (95.7)}}
$$

= 502 Btu/(hr) (sq.ft.) (9r)

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(b) h_o by Broaley Equation

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$$
h_0 = 0.728 \sqrt{\frac{k_f^3 \rho_f \rho_f \rho_r}{D_0 \mu_f \Delta t}}
$$

\n
$$
h_0 = 0.728 \sqrt{\frac{k_f^3 \rho_f \rho_f \rho_r \rho_s}{D_0 \mu_f \Delta t}}
$$

\n
$$
h_0 = 0.728 \sqrt{\frac{k_f^3 \rho_f \rho_f \rho_s}{D_0 \mu_f \Delta t}}
$$

\n
$$
= 512 \text{ Btu/(hr)} (sq, ft_s) (\Phi_f)
$$

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