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## The slow settling of a sphere in a viscous fluid in the proximity of a corner

Joseph Kisutcza New Jersey Institute of Technology

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#### THE SLOW SETTLING OF A SPHERE IN A VISCOUS

#### FLUID IN THE PROXIMITY OF A CORNER

BY

JOSEPH KISUTCZA

#### A THESIS

#### PRESENTED IN PARTIAL FULFILLMENT OF

#### THE REQUIREMENTS FOR THE DEGREE

OF

#### MASTER OF SCIENCE IN CHEMICAL ENGINEERING

AT

#### NEW JERSEY INSTITUTE OF TECHNOLOGY

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> Newark, New Jersey 1978

#### APPROVAL OF THESIS

THE SLOW SETTLING OF A SPHERE IN A VISCOUS

FLUID IN THE PROXIMITY OF A CORNER

BY

JOSEPH KISUTCZA

FOR

DEPARTMENT OF CHEMICAL ENGINEERING NEW JERSEY INSTITUTE OF TECHNOLOGY

BY

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APPROVED:

NEWARK, NEW JERSEY

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MAY, 1978

#### ABSTRACT

Experimental settling velocities for three different sizes of Delrin spheres in Ucon lubricant were determined at 20.2 degrees Celsius in order to confirm the validity of a theoretically derived equation for the settling of a sphere in the proximity of a corner. The experiments were conducted in a wedge shaped column with a circular sector base, filled with the viscous fluid, where the angle of the wedge was varied for experimental purposes. The distance from the wedge apex to the particle was also changed for the different runs.

The experimental data gave a good approximation of the values evaluated by the basic equation utilizing the drag force considering the wedge walls only. A modified form of the basic equation considering the additional drag from the vessel wall showed an improved agreement with the experimental data.

#### ACKNOWLEDGEMENTS

The author wishes to acknowledge the assistance and guidance given to him by Dr. E. Bart of this Institute during the course of this thesis.

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 $\sim 10^{-10}$ 

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 $\sim$ 

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 $\bar{\mathcal{A}}$ 



 $\sim$   $\sim$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$  ,  $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\bar{\omega}$ 

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#### CHAPTER I

#### INTRODUCTION

#### Scope and Purpose of Investigation

The topic of experimentally verifying the theoretical solutions for a sphere settling in the proximity of a corner in a viscous medium first come up in a discussion of its merits with E. Bart, author of a presently unpublished manuscript on the subject. In his work, Bart $^4$  derived expressions which mathematically evaluated the effects of two planes of arbitrary angles on a sphere settling parallel to their line of intersection. Investigation of the available literature showed that other authors $^{18,19}$  arrived at similar theoretical results, but experimental confirming data on the subject was non-existent.

These equations by Bart were derived for use in a medium bounded by an infinite wedge. Therefore a program was needed to evaluate the desired data where the theoretical conditions were approximated by actual equipment.

Aside from generating the experimental data, this author hoped that a modified form of Bart's equations for actual equipment might be established empirically if there was a considerable difference between the results of the experimental work and the calculated values from the basic equation. It should be noted that no rigorous solution applicable to a real container can be obtained,

but by piecing extant solutions together, a fairly accurate representation can be obtained.

It was anticipated that in order to achieve a good data fit, the effects of the vessel wall on the settling particle would have to be accounted for. Instead of deriving a new expression this author elected to augment Bart's work by a modified eccentricity and wall effect correction factor from the study of Greenstein and Happel $^{10}.$ 

The basic program for the investigation comprised the following:

- 1, to design and build an experimental system capable of simulating theoretical conditions and variables,
- 2, to gather experimental data on sphere settling in various wedge shaped domains filled with a viscous liquid,
- 3, to compare theoretical and experimental results and to evaluate the modified equations.

#### Literature Survey

The slow settling of particles in the presence of stationary surfaces has been of interest for many years. As early as 1896, Lorentz<sup>15</sup> treated the problem of a sphere slowly settling parallel to a plane wall by means of reflections. In 1907, Ladenburg<sup>14</sup> used the same technique to treat the settling of a sphere along the axis of an infinitely long cylinder at low Reynolds' number to a

first approximation. Since then, a multitude of solutions have appeared in the literature using the method of reflections to treat various configurations.

The problems concerning the slow motion of a sphere in the proximity of some stationary surface or surfaces can generally be broken down into three categories:

1, a particle moving parallel to a surface,

2, a particle moving either toward or away from a surface,

3, a combination of the two.

A solution presented by Sonshine, Cox and Brenner<sup>20</sup> for a sphere settling in a cylinder filled to a finite depth is an example of the last case. Any problem in which the sphere moves toward or away from a stationary surface must be of an unsteady nature. The literature is teeming with such solutions. Happel and Brenner $^{12}$ have reviewed many of the existing solutions. Since Happel and Brenner's overview was published, Sono and Hasimoto<sup>18,19</sup> have studied both categories 1 and 3 intensively.

Of greater concern here is the motion of a particle parallel to a flat surface or surfaces. These solutions can be of an unsteady nature if the particle is assumed to accelerate from rest. The more usual cases treat the steady motion of a particle parallel to a surface or surfaces. The aforementioned solutions of Lorentz $^{15}$ 

and Ladenburg<sup>14</sup> fall into this category. Faxen<sup>6,7,8</sup> investigated several problems in this category. He corrected Ladenburg's work for the cylinder and verified the solution of Lorentz for the sphere and flat plate. He has further extended Lorentz's problem by obtaining a torque solution and by obtaining higher ordered corrections. In addition, he solved the problem of a sphere settling parallel to and between parallel plates, of which the sphere and the flat plate problem is a special case. More recently Happel and Bart $^{11}$  treated a related problem of a sphere falling parallel to four walls, that is, the settling of a sphere along the axis of an infinitely long square duct. The solutions presented for a sphere settling in the proximity of a corner by Bart $^4$ , which formed the theoretical basis for this thesis appears identical to those derived by Sano and Hasimoto<sup>19</sup>. Although the presentation of the derivation was slightly different, these authors independently using the reflection method, arrived at the same end results. The expressions obtained by Bart seem to have a slight advantage over the work by Sano and Hasimoto since the former has obtained some higher ordered corrections for the translating particle and derived expressions for the rotational effects, while the latter considered only the first order effects without rotation.

Several additional investigators have examined var-

ious aspects of the problem of a sphere in a cylinder. Most notable, in terms of the present work, are those works concerning the settling of a sphere parallel to the axis at some distance. Happel and Brenner $^{12}$  presented a discussion showing how these solutions reduce in limiting cases to Lorentz's problem of a sphere and a flat plate.

Brenner and Happel<sup>5</sup> developed expressions and  $co$ efficients for the drag and torque for the translation of a single spherical particle in an infinitely long cylinder where the particle is kept from rotating. Later Greenstein and Happel $^{10}$  extended the problem treated by Brenner and Happel where the sphere may rotate and developed corrected values for the coefficients.

#### CHAPTER II

#### THEORY

#### Translation of a Sphere in a Corner

Consider a sphere oriented upon the midplane of a space formed by the intersection of two planes at some arbitrary dihedral angle. The wedge-shaped space thus formed is filled with an incompressible viscous liquid. The sphere is assumed to settle under the influence of gravity in a direction parallel to the apex of the wedge as shown in Figure 1. The angle of the wedge is arbitrary but must be sufficiently large so that the walls do not touch the sphere. Instead of using the wedge angle, it is more convenient to work with the half wedge angle,  $\phi_{\alpha}$ , since solutions symmetrical about the plane containing the sphere center are used. Thus, if  $\phi_{\alpha}$  is less than  $\pi/2$  the sphere is falling within a corner. If  $\phi_{0}$  is  $\pi/2$ , the sphere is falling parallel to a flat plane, which should yield values confirming the solutions of Lorentz $^{15}$ and Faxen<sup>6,7,8</sup>. When  $\phi_0$  is larger than  $\pi/2$ , the sphere is external to the corner and when  $\phi_{0}$  is equal to zero, a degenerate case occurs in which the wedge is a plane whose sharp edge faces the sphere. In the last case, the sphere settles in an otherwise unbounded fluid parallel to the sharp edge of an infinitely thin plate. The fluid velocity is zero upon both surfaces of this thin plate.







Figure 2 depicts the sphere-wall geometries described on the preceding page.

The equations to be solved are the creeping motion equation

$$
\mu \nabla^2 \overline{\mathbf{v}} = \nabla \mathbf{p} \tag{2.1}
$$

and the equation of continuity

$$
\nabla \cdot \overline{\mathbf{v}} = 0 \tag{2.2}
$$

The boundary conditions which define the fluid velocity are that

- 1, at the fluid-solid interface there is no relative motion;
- 2, the velocity at the sphere surface is the settling speed of the sphere.

The boundary value problem can be solved by a technique of successive approximations known as the method of reflections. For a comprehensive decription of the method, see the treatise on the subject by Happel and Brenner $^{12}.$ The solutions for fluid velocity, drag force, and torque may be obtained by summing the contributions of the individual fields.

 $\bar{\mathbf{v}} = \sum_{i=1}^{\infty} \bar{\mathbf{v}}^{(i)}.$ (2.3)  $\overline{F} = \sum_{i=-1}^{\infty} \overline{F}^{(i+2)}$ . (2.4)











b,  $\phi_0 = \pi/2$ 







 $d, \phi$  =  $\pi$ 

 $\hat{\mathcal{L}}$ 

 $\Delta_{\rm{eff}}$ 

$$
\overline{\mathbf{T}} = \sum_{i=-1}^{\infty} \overline{\mathbf{T}}^{(i+2)}.
$$
 (2.5)

To insure that the alternate velocity solutions are independent, the odd numbered solutions are unbounded at the sphere center and are zero infinitely far from the sphere, whereas the even numbered ones are finite at the sphere center and also vanish at a distance infinitely far removed from the sphere. This will insure that at large distances from the disturbing influence of the sphere the fluid velocity becomes zero. Because of this, only the odd numbered fields make contributions to the final drag and torque solutions.

When the results of the reflection solutions are summed in accordance with equations (2.4) and (2.5), the following expressions for the drag force and torque are obtained

$$
\vec{F} = 6\pi \mu \text{Ua} \vec{k} \left\{ 1 + f_1(\phi_0) (a/x_0) + f_1(\phi_0)^2 (a/x_0)^2 + [f_1(\phi_0)^3 + f_2(\phi_0)] (a/x_0)^3 \right\}
$$
\n(2.6)

and

$$
\bar{\mathbf{T}} = 4\pi \mu \text{Ua}^{2} \bar{\mathbf{j}} \{g_{1}(\phi_{0}) (a/x_{0})^{2} + f_{1}(\phi_{0}) g_{1}(\phi_{0}) (a/x_{0})^{3} + [f_{1}(\phi_{0})^{2} g_{1}(\phi_{0}) + g_{2}(\phi_{0}) (a/x_{0})^{4}].
$$
\n(2.7)

The power series may be replaced by the sum of a geometric progression, using techniques presented in Happel and Brenner<sup>12</sup>, to produce a still better approximations of the results for the drag force.

$$
\bar{F} = \frac{6\pi\mu\text{Ua}\bar{k}}{1 - f_1(\phi_0)(a/x_0) - f_2(\phi_0)(a/x_0)^3}.
$$
 (2.8)

A similar representation of the power series for the torque in powers of  $f_1(\phi_0)(a/x_0)$  would yield

$$
\bar{T} = 4\pi \mu U a^{2} \bar{j} \frac{q_{1}(\phi_{0}) (a/x_{0})^{2} + q_{2}(\phi_{0}) (a/x_{0})^{4}}{1 - f_{1}(\phi_{0}) (a/x_{0})}.
$$
 (2.9)

The settling particles ability to achieve free rotation will introduce some uncertainity in the solution for drag and torque. The coefficients of the powers of  $(a/x_{\Omega})$ are dependent of whether or not rotation is possible. This effect will not be obvious until  $\left( a/x_{\text{o}} \right)^4$  is reached in the drag solution and not until  $(a/x_{\overline{O}})^5$  is reached in the torque solution. Therefore, equations (2.6) and (2.7) are correct for both cases as far as the approximations are concerned.

## Evaluation of Terminal Settling and Angular Velocities in a Corner

The terminal velocity of a sphere settling in a corner may be evaluated from the drag obtained from either equation (2.6) or (2.8), depending on the degree of approximation desired and from Stokes law:

$$
\bar{F} = 6\pi \mu U_e a \bar{k} \tag{2.10}
$$

where

$$
U_{S} = \frac{2(\rho_{D} - \rho_{1})ga^{2}}{9\mu}.
$$
 (2.11)

Equating equation (2.10) and (2.6), it is apparent that

$$
U_{S}/U = 1 + f_{1}(\phi_{o}) (a/x_{o}) + f_{1}^{2}(\phi_{o}) (a/x_{o})^{2}
$$
  
+  $[f_{1}^{3}(\phi_{o}) + f_{2}(\phi_{o})] (a/x_{o})^{3}$ . (2.12)

Combining equation (2.8), using the geometric series approximation of higher ordered terms, with equation (2.10) yields

$$
U/U_{s} = 1 - f_{1}(\phi_{0}) (a/x_{0}) - f_{2}(\phi_{0}) (a/x_{0})^{3}
$$
. (2.13)

The angular velocity of the sphere as it falls will depend upon the spherical isotropy of the falling sphere. For a sphere where the centroid of the particle is not at the sphere center, the angular velocity must be zero. However, for a perfectly spherical freely rotating sphere, the torque necessary to prevent rotation must be

$$
\overline{T} = 8\pi\mu a^3\omega\overline{j}.
$$
 (2.14)

Combining this with either equation (2.7) or (2.9), yields, respectively

$$
\omega = U/2a\{g_1(\phi_0)(a/x_0)^2 + f_1(\phi_0)g_1(\phi_0)(a/x_0)^3
$$
  
+  $[f_1^2(\phi_0)g_1(\phi_0) + g_2(\phi_0)](a/x_0)^4$  (2.15)

and

$$
\omega = \frac{U/2 a [g_1 (\phi_0) (a/x_0)^2 + g_2 (\phi_0) (a/x_0)^4]}{1 - f_1 (\phi_0) (a/x_0)}.
$$
 (2.16)

The functions  $f_1(\phi_0)$ ,  $f_2(\phi_0)$ ,  $g_1(\phi_0)$  and  $g_2(\phi_0)$  have been evaluated numerically by Bart $^4$  for various values of

the parameter  $\phi_{\alpha}$  and the results are tabulated in Tables I and 2.

## Translation of a Sphere in a Cylindrical Tube

The inclusion of the sections on a sphere settling in a cylinder was necessitated by the fact that there are no infinite wedges in the real world. Therefore, to properly evaluate the experimental data, the wall effects must be evaluated and included in the final equation. The derivation of the equations dealing with the wall effects are shown in this and the following sections.

Consider the translation and rotation of a sphere moving with an arbitrary constant velocity through a viscous fluid in an infinitely long cylindrical tube. The sphere moves with a constant velocity parallel to the cylinder axis, displaced from the axis by some distance, as shown on Figure 3.

The fluid motion is governed by the creeping motion and continuity equations, (2.1) and (2.2), respectively. To solve these equations, the boundary conditions required are that

- 1, at the fluid-solid interface there is no relative motion,
- 2, at large distances from the disturbance caused by the moving sphere the velocity distribution becomes Poiseuillian.

The solution for the above problem makes use of the reflec-

## TABLE 1

VALUES OF  $f_1(\phi_0)$  AND  $f_2(\phi_0)$  FOR VARIOUS VALUES OF  $\phi_0$ 



 $\mathcal{F}^{\pm}$ 

 $\bar{z}$ 

# TABLE 2

VALUES OF  $g_1(\phi_0)$  AND  $g_2(\phi_0)$  FOR VARIOUS VALUES OF  $\phi_0$ 



 $\sim 10^{-1}$ 

## FIGURE 3

SPHERE SETTLING IN A CYLINDRICAL TUBE

 $\bar{\mathcal{A}}$ 



 $\sim$ 

tion method as previously described for a sphere settling in a corner.

The frictional force and the torque can be evaluated by adding the contributions of each field.

$$
\overline{\mathbf{F}} = \sum_{i=0}^{\infty} \overline{\mathbf{F}}^{(i)}.
$$
 (2.17)

$$
\overline{T} = \sum_{i=0}^{\infty} \overline{T}^{(i)}.
$$
 (2.18)

The final result for the frictional force for a sphere settling in a quiescent fluid where we set  $\beta = x_0/R_0$  is as follows:

$$
\bar{F} = 6\pi \mu \text{Ua} \bar{k} (1 + f(\beta) (a/R_0) + f^2(\beta) (a/R_0)^2)
$$
 (2.19)

and, for a freely rotating sphere, the torque is

$$
\bar{\mathbf{T}} = 8\pi \mu \text{Ua}^{2} \bar{\mathbf{j}} \{ g(\beta) (a/R_{0})^{2} [1 + g(\beta) (a/R_{0}) ] \}.
$$
 (2.20)

## Evaluation of Terminal Settling and Angular Velocities in a Cylinder

The terminal velocity of a sphere settling in a cylinder, offset from the cylinder axis, may be derived by combining equations (2.19) and (2.10):

$$
U/U_c = 1 - f(\beta) (a/R_c).
$$
 (2.21)

Similarly, for the angular velocity, equating (2.14) and (2.20) yields:

$$
\omega = U/a\{g(\beta)(a/R_0)^2[1+g(\beta)(a/R_0)]\}.
$$
 (2.22)

The functions  $f(\beta)$  and  $g(\beta)$  have been previously defined and reported in Happel and Brenner $^{12}$ , but an expanded and corrected set of values are presented in the work by Greenstein and Happel $^{10}$ . These latter values are listed in Tables 3 and 4.

#### Combined Equation for Wedge Contained in a Vessel

The equation (2.13) derived by Bart<sup>4</sup> accounts for the effect of the wedge on the settling particle, while the equation (2.21) presented by Greenstein and Happel $^{10}$ describes the wall effects. The mode of their derivation and the format, in which they are presented suggest that these equations may be combined to form an expression to estimate the settling velocities for a sphere in a column of viscous liquid, where the base of the column may be described as a sector of a circle with finite dimensions. In a column of a large diameter where the wall effects are small, as in the experimental vessel, this treatment should yield reasonably accurate results.

The proposed equation takes the form of Bart's equation (2.13) augmented by a modified term from equation (2.21). The modification consisted of reducing the calculated wall effect by a fraction which is the available circle segment divided by circle circumference. The final combined equation is as shown on the following page. 18

#### TABLE 3

TABULATION OF f(B) FOR VARIOUS VALUES OF B



 $\sim 10^{11}$ 

 $\label{eq:2.1} \frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^{2}}\left|\frac{d\mathbf{x}}{d\mathbf{x}}\right|^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d\mathbf{x}^{2}d$ 

### TABLE 4

TABULATION OF g(B) FOR VARIOUS VALUES OF B



 $\lambda$ 

$$
U/U_{s} = 1 - f_{1}(\phi_{0}) (a/x_{0}) - f_{2}(\phi_{0}) (a/x_{0})^{3}
$$
  
- f(\beta) (\phi\_{0}/\pi) (a/R\_{0}), (2.23)

where the coefficients  $f_1(\phi_0)$ ,  $f_2(\phi_0)$  and  $f(\beta)$  are as listed in Tables 1 and 3.

 $\label{eq:2.1} \frac{1}{2} \sum_{i=1}^n \frac{1}{2} \sum_{j=1}^n \frac{$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\label{eq:1} \nabla \mathbf{u} = \nabla \mathbf{u} + \nabla$ 

#### CHAPTER III

#### PHYSICAL PROPERTIES DETERMINATIONS

#### Fluid Medium Description

The fluid utilized in this experimental work was Ucon lubricant, type 50 HB-5100 from Union Carbide. It is a water soluble polyalkylene glycol type heat transfer and lubricating agent. Its stability under conditions encountered during testing, its density range, and its temperature-viscosity properties made it an excellent candidate for measuring slow settling velocities. Although the manufacturer's publication lists some of the desired physical properties, they were redetermined for the expected operating range.

#### Fluid Density Measurements

The density of the Ucon lubricant was determined by a modified version of ASTM Standard Test D-891 Method  $c^3.$ A brief description of the modified test procedure follows.

The bath was preset to the desired temperature and the calibrated 25 ml Gay-Lussac specific gravity bottle filled with Ucon lubricant was suspended in the bath. A second bottle equipped with a thermometer, also filled with the lubricant, was suspended next to the first bottle to check when thermal equilibrium was reached (usually 5-10 minutes). At the correct temperature the cover
was placed on the specific gravity bottle and the volume was adjusted. Upon removal from the bath the bottle was dried and weighed on a Satorius 3482/Electronic analytical balance and the data was recorded. The procedure was repeated for each desired temperature.

The experimentally generated data points were regressed linearly. The regression coefficients were calculated using the Curve Fitting Program  $SD-03A^{13}$  on a Hewlett-Packard 97 calculator. The correlation yielded the following equation:

 $p = 1.076666667 - 0.000758889T,$  (3.1) where  $\rho$  is the density in gms/cm<sup>3</sup> and T is the temperature in  $^{\circ}$ c.

The equation correlated to the data points appears to fit very closely, since the regression yielded a correlation coefficient of 0.999912117. The experimentally acquired fluid densities were plotted in Figure 4 as the function of the temperature.

The accuracy of equation (3.1) was checked by comparing the values generated by this equation to those listed by the manufacturer<sup>9</sup>. The agreement was excellent as shown in Table 5.

## Fluid Viscosity Measurements

The Ucon lubricant viscosity was evaluated according to ASTM Standard Test  $D-445^{\frac{1}{2}}$ . The experimental measure-

# FIGURE 4

# EXPERIMENTAL FLUID DENSITIES FOR UCON LUBRICANT



Temperature °C

# TABLE 5

 $\sim$ 

COMPARISON OF LIQUID DENSITIES FROM MANUFACTURER'S LITERATURE<sup>9</sup> WITH EQUATION (3.1) FOR UCON LUBRICANT



ments were carried out with a size 500 Cannon-Fenske viscometer. The viscometer was calibrated with water for an earlier unrelated experiment by the author. The viscometer calibration constant vs. temperature curve from the earlier work is reproduced in Figure 5.

The experimental viscosity data fitted to the type of equation developed by Watson, Wein and Murphy<sup>21</sup>. The regression coefficients were calculated on Hewlett-Packard 97 calculator using the Curve Fitting Program  $SD-03A^{13}$ . A variety of modifications to the basic equation was tried. The best results were achieved using the logarithmic curve fit which yielded the following:

$$
\mu = e^{\left\{e^{\left(5.138244744 - 0.561074273 \ln(1.8T-132)\right)} - 1.7\right\}},
$$
\n(3.2)

where  $\mu$  is the kinematic viscosity in centistokes and T is the temperature in <sup>O</sup>C.

The fitting of the equation to the data was very successful, since the regression produced a correlation coefficient of 0.999987224. Figure 6 displays the plot of experimental fluid viscosities vs. temperature.

The fluid viscosities derived by equation (3.2) were compared to the manufacturer's data<sup>9</sup>. This comparison is shown in Table 6. The agreement was very good, since the slight positive deviation may be explained by the low moisture levels (0.26%) in the lubricant. This phenomenon

# FIGURE 5

 $\mathcal{A}$ 

 $\bar{\mathcal{A}}$ 

 $\bar{z}$ 

 $\mathcal{A}^{\mathcal{A}}$ 





Temperature °C

 $\sim$ 

# FIGURE 6

EXPERIMENTAL KINEMATIC VISCOSITIES FOR UCON LUBRICANT



 $\epsilon$ 

## TABLE 6

# COMPARISON OF LIQUID VISCOSITIES FROM MANUFACTURER'S LITERATURE<sup>9</sup> WITH EQUATION (3.2) FOR UCON LUBRICANT



 $\mathcal{L}^{\pm}$ 

of a small increase in viscosity of fluids of this type at low levels of contained water is documented in the manufacturer's literature<sup>9</sup>.

### Sphere Description

The spheres used for the experimental work were Delrin spheres of Grade 200 from Ultraspherics. The polymer used in manufacturing the spheres was developed by E.I. DuPont. Delrin is an opaque white, acetal type polymer. These spheres are normally used in highly critical bearing applications and they are highly polished. These spheres were selected for their stability under normal experimental conditions and for their relative density to the fluid medium. The sphere sizes acquired were 5/32, 1/4 and 11/32 inch nominal diameters.

## Sphere Selection Process

The applications for which these spheres were designed, required that their basic diameter tolerance be very low: therefore, sizewise, they are nearly identical. However, spotchecking revealed that there was a considerable variation in densities for the same sizes and even a larger difference was found between the different ones.

A procedure was instituted to select a number of spheres of each size with uniform densities. An abbreviated account of the procedure is listed below.

A solution of 350 gms of Tetrachloromethane (MCB

Spectroquality, S.G.=1.5940) and 1, 2-Dichloroethane (MCB Spectroquality, S.G.=1.2351) was placed in a 6 inch diameter glass cylinder and 100 each of the spheres of 5/32, 1/4 and 11/32 inch nominal diameter were placed in the solution. All the spheres sunk to the bottom of the cylinder. Tetrachloromethane was added to the solution at 1.25 ml increments and the resulting solution was stirred. After stirring the solution was allowed to come to rest. Prior to each addition all spheres that have risen from the bottom were collected and segregated by size and approximate density. When large segments of the spheres of each size were collected, the density of the solution was also determined.

For each nominal diameter, the group with the largest number of spheres with the same approximate density was selected. Each sphere from these groups was weighed individually and was subjected to a multiple point determination of its diameter. Of the ones which appeared identical, six were selected at random for determination of their exact densities.

### Sphere Density Determination

The exact densities for the spheres were determined by a modified version of ASTM Standard Test D-167 $^2$ . A brief description of the procedure used follows.

A 25 ml Walker type specific gravity bottle was calib-

rated. The six spheres were weighed collectively and placed in the bottle. The bottle was filled with 1,2-Dichloroethane (MCB Spectroquality) and placed in the constant temperature bath. When the solvent reached thermal equilibrium at 20.0 degrees Celsius, the volume was adjusted. From the resulting measured volumes, densities for the spheres were calculated. As a check on the measured diameters, the sphere volumes were also used to obtain calculated diameters.

The resultant properties for the spheres are tabulated in Table 7.

# TABLE 7

### SELECTED DELRIN SPHERE DIAMETERS AND DENSITIES



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#### CHAPTER IV

#### EXPERIMENTAL SYSTEM

### Design Considerations

The following considerations influenced the overall design of the experimental system:

- 1, a need for the largest possible diameter vessel to minimize the wall effects, but where the contained liquid is still transparent to allow a clear view of the settling particle;
- 2, a need for a stable platform to provide support for the sphere release mechanism and for the wedges suspended in the liquid;
- 3, a need for all internal parts to be constructed from translucent materials;
- 4, a need for a means to recover the spheres from the bottom of the tank without greatly disturbing the system;
- 5, a need for a constant temperature environment to minimize the temperature fluctuations in the liquid;

### Experimental Equipment

Since none of the available equipment fitted the above mentioned considerations, a decision was made to design and build the equipment for use specifically in this

study. Although the overall design is unique, the design took advantage of commercially available pieces of equipment wherever possible. These pieces, with minor modifications, became part of the overall design. Sections, which were radically different from existing equipment, were designed for ease of use and with minimum expenditure of materials. The equipment, as designed, consisted of a large container with a dual purpose support and alignment platform, a sphere release mechanism and four wedges of various angles. The schematic of the equipment as built is shown on Figure 7.

### Ucon Lubricant Container

The preliminary experiments for the determination of the maximum width of the liquid which does not impair the observation of the settling particle showed that when the viewing path exceeds 30 inches the observation becomes difficult. The optimum width of the viewing path through the liquid was found to be between 25 and 30 inches, where the Ucon lubricant takes on a deep green hue but stays transparent, therefore a 24.0 inch (I.D.) by 36.0 inch (T.L. to T.L.) by 3/32 inch (wall thickness) vessel was selected as the basic container for the sphere settling experiments. The container, prior to modifications, was an open top, dished bottom head feedtank of 316 S.S. construction with a 1 inch bottom drain. The container

## FIGURE 7

# SCHEMATIC OF EXPERIMENTAL EQUIPMENT



was supported on 3 tubular legs with casters and a bolt type levelling assembly on each leg.

To comply with the desired design basis, various modifications were installed on the basic container. The description of these alterations are listed below.

1, Four custom made viewing ports were constructed from 4 inch I.D. 316 S.S. tube stub ends by placing a 4.75 inch diameter by 1/4 inch port glass, protected on both sides with CRT envelope gaskets, between a retaining ring and the flat of the stub end. The retaining ring and the flat of the stub end were drilled out in four places and bolted together. The container had four 4 inch holes (2 on each side) cut on 12 and 24 inch centers from the bottom tangent line and each of the assembled viewing ports were seal welded to the container. The internal weldseam and any other protrusions were ground to a mill finish.

2, A 5.5 inch high by 6 inch top radius half round powder funnel of 316 S.S. construction was force fitted into the bottom drain coupling with the round part toward one set of viewing ports. The segment of the tank which contained the circular portion of the funnel was designated as the front.

3, Two 1/8 inch compression fittings were attached to the vessel to act as the thermocouple connections. They were located 1 inch below the top tangent line

and 1 inch above the bottom tangent line on the same vertical as the rear viewing ports.

4, An isolation reservoir, similar to the one employed by Matyas $^{16}$  in his experimental work, was constructed from two 1 inch ball valves and a 1.5 inch I.D. by 3.5 inch long sight glass. The pieces were connected together with 1 inch minimum length pipe nipples. An identical nipple was used to join the completed isolation reservoir to the bottom coupling.

Figure 8 is a sketch of the Ucon lubricant container as used in the experiments, showing some of its critical dimensions.

# Wedge Support and Alignment Platform with Sphere Release Mechanism Alignment Assembly

A platform, as shown on Figure 9, was designed for dual purpose and was constructed from 3/4 inch by 3/16 inch 316 S.S. barstock. Five of the arms are single layer construction and each of the arms had a hole drilled 1 inch from the outer end. These arms supported and aligned the wedges in the liquid. The three other arms of double construction had a second bar attached 0.6 inches above the lower ones. The double arms had 5 holes drilled through both bars at 1 inch intervals from the center. These double arms supplied the vertical and radial alignment for the sphere release mechanism. The platform was also



**FIGURE**  $\infty$ 

UCON LUBRICANT CONTAINER

 $\star$ 

WEDGE SUPPORT AND ALIGNMENT PLATFORM WITH SPHERE RELEASE MECHANISM ALIGNMENT ASSEMBLY



drilled out at its center for the central wedge support. All holes on the platform were drilled with a No. 12 drill bit (0.1890 inch I.D.).

#### Sphere Release Mechanism

A Triceps type forceps Model T8 was modified to handle the positioning and release of spheres in the liquid. The modification consisted of attaching a 3/4 inch washer to the forceps with epoxy cement 1.5 inches from the top to act as a stop for its vertical travel. This unit inserted through the proper hole on the double arm gives a stable and reproducible starting point for the sphere during the experimental runs.

Figure 10 shows the sphere release mechanism.

### Wedge Sections

The wedge sections were fabricated from Plexiglas brand 1/8 inch thick acrylic sheet (ANSI Z97.1-1966/72 079U) from Rohm & Haas. Four wedges were produced, each forming a different angle (60, 90 which also doubled as the 270, 180 and 360). Each wedge was 34 inches high with the sides having a radial distance of 12 inches. The support rod and the brackets were formed from 316 S.S. 10-24 threaded rods and 5/8 inch by 1/8 inch 316 S.S. channels respectively. The threaded rods were spot welded to the top of the channel. For the angled brackets the channels were cut and welded together to form the correct

# FIGURE 10





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 $\overline{\phantom{a}}$ 

angle before the threaded rod was attached. The completed brackets were placed at predetermined locations on the top edge of the wedge and each had two 1/8 inch holes drilled through both the channel and the acrylic sheet. The brackets were attached to the wedges through the predrilled holes with short 1/8 inch sheet metal screws.

The 180 degree wedge was formed from two 34 inch by 12 inch sheets, which were connected together with small hinges near the top and bottom edges. The hinging allowed the insertion of the wedge past the restriction on the top of the lubricant container. The central support on the 180 degree wedge was connected in place after the insertion into the container.

Figures 11 thru 14 are the assembly drawings of the wedge sections for the various angles.

### Auxilary Equipment

The temperature of the lubricant was constantly monitored during the experiments at 1 inch below the liquid surface and at the bottom of the container. The temperature measurement was accomplished by the use of calibrated 1/16 inch Chromel-Alumel thermocouples. Each of the thermocouples were connected to a CONDEC digital indicator, which provided continuous readout of the temperatures. The range of the instrument was 999.9 degrees Celsius with 0.1 degree accuracy.

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ASSEMBLY DRAWING FOR THE 60° WEDGE SECTION





ASSEMBLY DRAWING FOR THE 90<sup>°</sup> WEDGE SECTION



Bevel mating edges to 45° and attach



ASSEMBLY DRAWING FOR THE 180° WEDGE SECTION





ASSEMBLY DRAWING FOR THE 360° WEDGE SECTION

 $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\bar{z}$ 

 $\mathcal{L}$ 

FIGURE 14

For the determination of the settling times, a Faehr brand digital stopwatch was utilized. The timer has a 5 digit display capability with 0.1 second accuracy.

The background lighting for the container was provided by a 2 feet long 32 watt fluorescent light, which was placed approximately 15 inches beyond the rear viewing ports.

### Testing Environment

To overcome the temperature fluctuation which plagued earlier experimental work $^{16}$ , the experimental equipment was placed in an 8 feet high by 7 feet by 6 feet Geldback temperature controlled enclosure, where all the sphere-dropping experiments were conducted.

To insure that the system reached thermal equilibrium, the control unit on the enclosure was set to 68 degrees Fahrenheit three days prior to the start of the experiments. During these days the lubricant and the enclosure air temperature was monitored.

Within 3 hours the lubricant temperature reached equilibrium at 20.2 degrees Celsius and remained there without changing. The equilibrium air temperature inside the enclosure measured 68.4 degrees Fahrenheit. Opening the door on the enclosure changed the temperature less than 1 degree Fahrenheit and after the door was closed the temperature returned to equilibrium within 3 minutes. Short openings of the door did not effect the lubricant temperature.

The enclosure was located within a larger room where the temperature was kept between 66 and 72 degrees Fahrenheit, which also helped to stabilize the temperature fluctuation in the controlled enclosure. This location was also used to store all equipment not in use.

#### Test Conditions and Limitations

To test the validity of the derived equations, all combinations of variables used in the equations were tested, with the exception of those which were the functions of the temperature. The lubricant temperature was kept constant at 20.2 degrees Celsius. In all cases, sextuplicate runs were made to test reproducibility. Agreement between the six runs never varied more than 3 percent.

As was explained earlier, that the primary limitation was the size of the lubricant container, which further restricated some of the other variables. The experiment was designed so that the wedge apex to particle center distance should vary up to approximately 40 percent of the container radius to reduce the effects of the container wall on the setting particle.

The number of different particle sizes were limited by the availability of various size spheres of proper

grade and material. Since these precision spheres are normally custom manufactured, we were fortunate to acquire a good selection of each of 3 widely diverse sizes. The larger sizes were approximately 1.6 and 2.2 times the diameter of the smallest one.

The final variable, the wedge angle had limitations imposed on by the theoretical work on which this experiment was based. The coefficients used in the equations had been calculated for only a few selected angles; therefore, the comparison of experimental to calculated results would not have been possible even if there were more wedges built. Only for one angle was data gathered where there were no coefficients calculated since the derivation of coefficients for the 270 degree angle was underway when the experimentation began, although at the writing of this thesis, it is still not completed.

#### Testing Procedure

A detailed description of the experimental procedure is listed below.

For each wedge angle the listed procedure was followed:

1, The proper wedge was lowered into the liquid and the support rods were attached loosely to the wedge support platform.

2, The container walls were tested for verticality with a long bubble type carpenters level. If it was needed, the leveling was accomplished by adjusting

the bolts on the leveling assemblies on each leg. 3, The wedge support platform was rotated to the proper orientation, i.e., the sphere release mechanism alignment bar was turned to the viewing port to viewing port axis.

4, The top edge of the wedge was leveled by the aid of the bubble level. The adjustment was done by tightening or loosening the nuts on the support rods. 5, After all disturbance of the liquid ceased, the system was allowed to come to equilibrium for a minimun of a half hour.

For each distance from the wedge apex to the particle center the procedure was as follows:

1, The sphere release mechanism was inserted through both holes at appropriate locations on the alignment bar.

2, All spheres were wetted with Ucon lubricant.

3, The sphere release mechanism was raised and a sphere was placed in the clampa.

4, The release mechanism was lowered into the liquid until the stop on it impeded the downward travel. 5, The sphere was released by pressing down the plunger on the top of the release mechanism and holding it down for 10 seconds.

6, The timer was started when the sphere interected the plane formed by the timing marks on the upper viewing ports and was stopped when the sphere reached the plane formed by the timing marks on the lower viewing ports.

For each sphere diameter and repetitions, Steps 3 thru 6 of this section were repeated.

7, After 18 spheres were dropped, the upper valve of the lock system was closed and the lower one opened, thereby draining out the Ucon lubricant and the spheres.

8, When all spheres were removed, the lower valve was closed and the upper one opened and the lock was allowed to fill with lubricant again. The removed spheres were readied for other runs by draining the excess lubricant back to the container.

9, After all disturbance of the liquid ceased, the system was allowed to come to equilibrium for a minimum of a half hour.

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#### CHAPTER V

#### EXPERIMENTAL RESULTS

#### Analysis of the Results

The experimentally determined settling times were converted to experimental settling velocities by taking the reciprocal of the settling times. This simple conversion was made possible since the distance where the settling was measured was exactly one foot. For each combination of variables, six determinations were made; therefore a mean value for the settling velocity and a standard deviation for the set were calculated. To examine the data scatter a conversion of the calculated standard deviations was required. The format where the comparison gave meaningful results was arrived at by dividing the standard deviation for each set by the calculated mean settling velocity for the same set. This data,  $\sigma/U_{\text{em}}$ , was plotted in Figures 15, 16 and 17 for each sphere size as the function of the distance the particle is from the wedge apex to the vessel radius ratio  $(x_0/R_0)$ . To determine if the plotted data followed a trend, the maximum and minimum values for  $\sigma/U_{em}$  for each  $x_0/R_o$  was used to obtain two lines representing the approximate limits for the data scatter by regressing the above mentioned data by the least squares method. These lines are also displayed on Figures 15, 16 and 17.

# FIGURE 15

# DATA SCATTER VARIATION FOR THE 5/32" SPHERE WITH THE DISTANCE FROM THE WEDGE APEX







DATA SCATTER VARIATION FOR THE 1/4" SPHERE WITH THE DISTANCE FROM THE WEDGE APEX





 $\mathcal{A}^{\mathcal{A}}$  and  $\mathcal{A}^{\mathcal{A}}$  and  $\mathcal{A}^{\mathcal{A}}$ 

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# FIGURE 17

# DATA SCATTER VARIATION FOR THE 11/32" SPHERE WITH THE DISTANCE FROM THE WEDGE APEX

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The slopes of the lines representing the approximate limits of the data indicate an inversely proportional relationship between  $\sigma/U_{\text{cm}}$  and  $x_{\text{o}}/R_{\text{o}}$ . The reason for this effect could be traced back to the rotation of the particle imparted by the walls of the wedge. The effect of the rotation of the sphere was impossible to determine in the equipment used for this work; therefore its effect is not included in the final equations.

Although extreme care was exercised in the selection of the spheres, there was a definite possibility of the particles having non-uniform internal densities, i.e., the centroid of the particle is not at the sphere center. With the existence of spheres rotating with various angular velocities, the likelihood of a whole spectrum of solutions is possible.

#### Comparison of Theory (Equation (2.13)) and Experiment

A dimensionless form for the observed and predicted sphere settling velocities  $(U/U_g)$  in Ucon lubricant at 20.2 degrees Celsius, as determined by the experiment and equation (2.13) for wedge angles of 60, 90, 180 and 360 degrees, are shown in graphical form as the function of x o /a in Figures 18 thru 21, respectively. This equation predicted the settling velocities based on the effects of the wedge wall. In all cases, the experiment gave good agreement for angles less than 180 degrees. For the 360 degree angle, the experimental values appear to deviate



 $\langle\sigma\rangle_{\rm eff}$ 

5/32" sphere  $\rightarrow$  $1/4"$  sphere of this experiment 11/32" sphere --- Equation (2.13)

COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION  $(2.13)$  FOR THE  $60^{\circ}$  WEDGE ANGLE

FIGURE 18

 $\Delta$
# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION  $(2.13)$  FOR THE 90<sup>°</sup> WEDGE ANGLE



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# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION  $(2.13)$  FOR THE 180<sup>°</sup> WEDGE ANGLE





# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION  $(2.13)$  FOR THE 360<sup>°</sup> WEDGE ANGLE





from the theory by a small margin. This outcome may be diagnosed by reviewing the physical make-up of the equipment. The calculations in the theory were based on a sphere settling parallel to an infinitely thin plate, whereas in the actual test equipment a plate with 1/8 inch thickness was substituted. The additional drag from the edge could be the cause for the observed deviation.

#### Comparison of Theory (Equation (2.23)] and Experiment

The modified equation (2.23), accounting for both the wedge and vessel wall effects, was compared to the experimental data and shown graphically in Figures 22 thru 25 for the wedge angles 60, 90, 180 and 360 degrees. The agreement of the predicted dimensionless form of the settling velocities with the experimental data is excellent for all cases, showing a considerable improvement over equation (2.13). The largest deviation from the predicted values was observed for the 360 degree angle. The explanation offered for this deviation is identical to the one proposed in the preceding section for equation (2.13). The agreement for the experimental and calculated values are particularly striking for angles less than 180 degrees where the differences are less than 1 percent for most cases.

Empirical Determination of Coefficients  $f_1(\phi_0)$  and  $f_2(\phi_0)$ At the writing of this thesis the coefficients

# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION  $(2.23)$  FOR THE  $60^{\circ}$  WEDGE ANGLE





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# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION  $(2.23)$  FOR THE 90<sup>°</sup> WEDGE ANGLE





 $\mathfrak{b}4$ 

# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION (2.23) FOR THE 180<sup>°</sup> WEDGE ANGLE





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 $\sim 10^7$ 

# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION (2.23) FOR THE 360<sup>°</sup> WEDGE ANGLE

 $\sim$ 





 $x_0/a$ 

 $f_1(\phi_0)$  and  $f_2(\phi_0)$  for the 270 degree angle had still not been derived; therefore, the experimental settling velocities were used to obtain these coefficients. Equation (2.23) was converted to the following format:

$$
-(a/x_0) f_1(\phi_0) - (a/x_0)^3 f_2(\phi_0) =
$$
  
(U/U<sub>S</sub>) + {f(\beta) (\phi\_0/\pi) (a/R<sub>O</sub>)} - 1, (5.1)

in order to facilitate the calculation of the coefficients by determinants. Using Cramer's Rule, the following equations were derived for the coefficients:

$$
f_1(\phi_0) = (b_i a_{j2} - b_j a_{i2}) / (a_{i1} a_{j2} - a_{j1} a_{i2})
$$
 (5.2)

and

 $f_2(\phi_0) = (b_i a_{i1} - b_i a_{i1}) / (a_{i1} a_{i2} - a_{i1} a_{i2})$ (5.3) where  $b_i$  and  $b_j$  is  $(U/U_s) + \{f(\beta) (\phi_0 / \pi) (a/R_o) \} - 1$ ,  $a_{11}$  and  $a_{j1}$  is  $-(a/x_0)$ , and  $a_{i2}$  and  $a_{j2}$  is  $-(a/x_0)^3$ , with i and j referring to two different linear equations for the same wedge angle.

All combinations of the available data were evaluated and an arithmetic mean and the deviation from the mean were determined.

The mean values for the coefficients determined by the above method are tabulated in Table 8.

The empirically determined coefficients were substituted into equations (2.13) and (2.23) and the calculated settling velocities, determined by this method,

## TABLE 8

# EMPIRICALLY DERIVED COEFFICIENTS FOR THE 270 DEGREE WEDGE ANGLE



 $\mathbb{Z}^2$ 

compared to the experimental values. The results of this comparison are displayed on Figures 26 and 27.

 $\bar{z}$ 

# COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION (2.13) USING THE EMPIRICALLY DERIVED COEFFICIENTS FOR THE 270<sup>°</sup> WEDGE ANGLE





 $\overline{c}$ 

COMPARISON OF EXPERIMENTAL SETTLING VELOCITIES WITH THE THEORETICAL PREDICTION OF EQUATION (2.23) USING THE EMPIRICALLY DERIVED COEFFICIENTS FOR THE 270 $^{\circ}$  WEDGE ANGLE

> $\triangle$  5/32" sphere 1  $\overline{1/4}$ " sphere of this experiment  $-$  11/32" sphere  $--$  Equation (2.23)



 $L<sub>1</sub>$ 

#### CHAPTER VI

#### CONCLUSIONS AND RECOMMENDATIONS

#### Conclusions

The results of the experiments conducted in the course of this thesis to determine the settling velocities of Delrin Spheres in Ucon lubricant may be summarized as follows:

- 1, Equation (2.13), using only the correction from the wedge wall effects, gives a reasonable approximation for the velocity of a particle settling in a wedge shaped domain of viscous liquid.
- 2, The modified equation (2.23), combining the vessel wall effects and the previously utilized wedge correction, accurately describes the settling of a sphere in a column of viscous fluid with a base of a circular sector.
- 3, The wedge walls impart a rotational effect of varying degrees on the particle settling near them. This induces a slight effect upon the translational velocity thereby increasing the scattering of the measured data. The unpredictable rotation is caused by the particles having non-homogeneous internal densities or by not being perfectly spherical.

4, The vessel wall contributes only a small portion

of the drag on the settling particle and it may be approximated by taking a fraction of equation (2.21).

5, The empirically derived coefficients for the 360 degree wedge angle differ only by a small amount from the theoretical values.

#### Recommendations

- 1, A study, effectively the continuation of this one, would be to determine the settling velocities near the vessel wall to test the validity of the derived equation for that region and expand its range. Experiments with other fluids and sphere materials could be used to further increase the confidence in the equation.
- 2, The rotation of the particle should be studied experimentally to quantify its effect on a settling particle, if spheres with uniform internal densities can be found.
- 3, A study not developed here would be the investigation of the effect the bottom of the vessel has on the particle settling toward it.
- 4, A study of particular importance would be that of the effect of the settling of multiple particles, which would extend the effectiveness of the equation for use in designing practical equipment.

## APPENDIX A

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#### SAMPLE CALCULATION

The physical properties determined prior to the start of the experimental runs and the settling times taken during a run were required to ascertain the following:

- 1, experimental settling velocities,
- 2, calculated settling velocities considering only the effects of the wedge on the settling particle,
- 3, calculated settling velocities based on the effects of both the wedge and the vessel wall.

A typical calculation for the above quantities is listed in this section. The data used for the sample calculation was collected for the run, where the wedge angle was 60 degrees and a 0.1562" diameter sphere was dropped 1 inch from the wedge apex.

#### Experimental Settling Velocities

The experimental settling velocities were calculated by taking the reciprocal of the recorded settling times, since the distance where the settling was measured was 1 foot.

 $U_{e(i)} = 1/\tau_{(i)}$ 

where  $U_{e(i)}$  is the experimental settling velocity in ft/sec and  $\tau$ <sub>(i)</sub> is the measured experimental settling time in seconds, with i referring to the replicate run for the same set.

Input data

 $\tau_1 = 362.0 \text{ sec.}$  $\tau_2$  = 363.4 sec.  $\tau_3 = 364.5 \text{ sec.}$  $\tau_4$  = 363.7 sec.  $\tau_5 = 362.6 \text{ sec.}$  $\tau_6 = 363.6 \text{ sec.}$ 

Results

 $U_{el}$  = 0.002762 ft/sec.  $U_{e2} = 0.002752$  ft/sec.  $U_{eq} = 0.002743$  ft/sec.  $U_{e4} = 0.002750$  ft/sec.  $U_{e5} = 0.002758$  ft/sec.  $U_{e6}$  = 0.002750 ft/sec.

An arithmetic mean for the set was calculated. This value was compared to the theoretical predictions.

$$
U_{em} = \left(\begin{array}{c} n \\ \sum U_{e(i)} \end{array}\right) / n
$$

where  $U_{em}$  is the mean experimental settling velocity in ft/sec and n is the number of replications.

Input data

The above calculated experimental settling velocities.

Result

U em = 0.002753 ft/sec.

The standard deviation for the set was also determined to evaluate the data scatter.

$$
\sigma = \left(\sum_{i=1}^{n} (U_{e(i)} - U_{em})^2 / n\right)^{\frac{1}{2}}
$$

where  $\sigma$  is the standard deviation from the mean for the set.

#### Input data

The experimental and experimental mean settling velocities listed on the previous page.

Result

 $\sigma = 0.6096 \times 10^{-5}$ 

## Calculated Settling Velocities

Since the calculated settling velocities are based on the Stokes settling velocity, therefore a value for the latter was evaluated.

 $U_{s} = [2ga^{2}(\rho_{p}-\rho_{1})]/9\mu$ 

where g is the acceleration of gravity in  $ft/sec^2$ , a is the particle radius in ft,  $\rho_p$  and  $\rho_1$  are the densities for the particle and liquid, respectively in  $1b/ft^3$ , and p is the liquid viscosity in lb/ft-sec.

Input data

 $g = 32.2 \text{ ft/sec}^2$ .

a = 0.0781 inch  
\n= (0.0781 inch) (0.0833 ft/inch)  
\n= 0.0065 ft.  
\n
$$
\rho_p
$$
 = 1.3883 gms/cm<sup>3</sup> from Table 7.  
\n= (1.3883 gms/cm<sup>3</sup>) (62.4264 cm<sup>3</sup>·lb/gms·ft<sup>3</sup>)  
\n= 86.6666 lb/ft<sup>3</sup>.  
\n $\rho_1$  = 1.0613 gms/cm<sup>3</sup> from Equation (3.1) for 20.2<sup>o</sup>C.  
\n= (1.0613 gms/cm<sup>3</sup>) (62.4264 cm<sup>3</sup>·lb/gms·ft<sup>3</sup>)  
\n= 66.2531 lb/ft<sup>3</sup>.  
\n $\mu$  = 2706.57 centistokes from Equation (3.2) for 20.2<sup>o</sup>C.  
\n= (2706.57 centistokes) (1.0764 x 10<sup>-5</sup> ft<sup>2</sup>/centi-  
\nstokes: sec) (66.2531 lb/ft<sup>3</sup>)  
\n= 1.9302 lb/ft/sec.

Result

 $U_{S} = 0.003206 \text{ ft/sec.}$ 

The calculated settling velocity considering only the wedge effects on the settling particle was evaluated.

 ${\tt U}_1\,=\, {\tt U}_{{\tt s}}[ \,1\!\!-\!(a/x_{{\tt o}})\, {\tt f}_1\,(\phi_{{\tt o}})\!\!-\!(a/x_{{\tt o}})\, ^3 \!{\tt f}_2\,(\phi_{{\tt o}})]$ 

where  $U_1$  is the calculated settling velocity in ft/sec, x o is the distance from the wedge apex to the particle center in ft, and  $f_1(\phi_0)$  and  $f_2(\phi_0)$  are the wedge angle coefficients for translating particles.

Input data



 $f_1(\phi_0) = 1.7891$  $f_2(\phi_0) = -2.7820$ U<sub>S</sub> = 0.003206 ft/sec.

Result

 $U_1 = 0.002761$  ft/sec.

Similarly, the calculated settling velocity considering the wedge and vessel wall effects on the settling particle was evaluated.

$$
U_2 = U_S [1 - (a/x_0) f_1 (\phi_0) - (a/x_0)^3 f_2 (\phi_0) - (a/R_0) (\phi/180) f(\beta)]
$$

where  $U_2$  is the calculated settling velocity in ft/sec,  $R_{\odot}$  is the fluid container radius in ft,  $\phi$  is half of the wedge angle in degrees, and  $f(\beta)$  is the eccentricity coefficient  $(x_0/R_0)$ .

Input data

 $= 0.0065$  ft.  $\mathbf{a}$  $x_0 = 0.0833$  ft.  $f_1(\phi_0) = 1.7891$  $f_2(\phi_0) = -2.7820$  $= 1.0$  ft.  $R_{\Omega}$ 

= 30 degrees  $\phi$ 

> An interpolated value was used for  $f(\beta)$ . From Table 4 the following was acquired:  $f(\beta)=2.10270$  for  $\beta=0.05$ and  $f(\beta)=2.09758$  for  $\beta=0.10$ . The interpolation

yielded  $f(\beta)=2.0993$  for  $\beta=0.0833$ .

 $f(\beta) = 2.0993$ 

Result

 $U_2 = 0.002754 \text{ ft/sec.}$ 

APPENDIX B

## TABLE 9

PHYSICAL PROPERTIES FOR UCON LUBRICANT 50-HB-5100<sup>9</sup>



 $\sim$   $\sim$ 

 $\mathcal{A}$ 

 $\sim$   $\mu$  .

APPENDIX C

 $\mathcal{A}$ 

## TABLE 10

 $\sim$   $\sim$ 

 $\hat{u}$ 

# PHYSICAL PROPERTIES FOR ACETAL (DELRIN) SPHERES $^{17}$



 $\epsilon$ 

APPENDIX D

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\right)\frac{1}{\sqrt{2}}\right)\frac{1}{\sqrt{2}}\,d\mu$ 

 $\theta$  .

### SOURCE LISTING FOR COMPUTER CALCULATIONS

The repetitive calculations to process the accumulated data were executed on Univac Series 70 computer in the N. J. I. T. Computer Center.

The source listing for the calculation and printing routine is included in this section.

The explanation of the special nomenclature used in the program is included in the source listing.

**1 INTEGER TBL1(5),TBL2(5,5),TBL3(5) 2 REAL TEMP(5,3,5),RHOL(5,3,5),MUKL(5,3,5),uS(5,3,5), 3 1MUAL(5,3,5),U1(5,3,5),U2(5,3,5),FBETA(5,3,5),TAU 4 2(5,3,5,6),UEXP(5,3,5,6),UEXPAv(5,3,5),SIGMA(5,3,5), 5 3PHIX2(5) 6 REAL\*4 DIAP(3)/0,1562,0.2497,0.3435/,RHOP(3)/1.3883, 7 11,3774,1,4001/,FPH11(5/1,7891,1,1584,0,5625,'N/A',**  8 20.4775/,FPH[2(5)/-2.7820,.0.8416,-0.125,'N/A', **9 3-0.05305/,X0(5)/1.0,2,0,3,0,4,0,5.0/,PHI(5)/30.0, 10 445.0,90.0,135.0,180.0/;RC/12.0/,G/32.2/,TEST/'N/A'/**  ii **C 12 C DIAP IS THE CALCULATED AND MEASURED DIAMETER OF THE 13 C PARTICLE IN INCHES 14 C RHOP IS THE CALCULATED DENSITY OF THE PARTICLE IN GMS/CC 15** C FPHI1 IS **THE FIRST COEFFICIENT IN THE WEDGE CORRECTION 16 C EQUATION DIMENSIONLESS i7 C FPHI2 IS THE SECOND COEFFICIENT IN THE WEDGE CORRECTION] 18 C EQUATION DIMENSIONLESS 19 C X0 IS THE DISTANCE FRCM THE WEDGE APEX TO THE PARTICLE 20 C CENTER IN INCHES 21 C PHI IS 1/2 OF THE WEDGE ANGLE IN DEGREES 22 C RO IS THE TANK RADIUS IN INCHES 23 C 0 IS THE ACCELERATION OF GRAVITY IN FT/SEC\*\*2 24 C TEST CHECKS FOR THE AVAILABILITY OF COEFFICIENTS IN THE 25 C EQUATIONS 26 C 27 REAL\*4 BETAL(26)/0.0.0.01,0.02,0.03,0.05,0.10,0.15,**  <sup>28</sup>**10.20,0.25,0.30,0.35,0,37,0.39,0.40,0.41,0.43,0,45, 29 20.50,0.55,0.60,0.65,0.70,0.75,0.80,0.85,0.90/ 30 C 31 C PETAL IS THE RATIO OF THE DISTANCE FROM WEDGE APEX TO THE 32 C PARTICLE CENTER OVER THE TANK RADIUS DIMENSIONLESS 33 C 34 REAL\*4 FBETAL(26)/2.10444,2.10433,2.10415,2.10381, 35 12.10270,2.109758,2.08962,2.07937,2.06801,2.05687, 36 22.04800,2.04561,2.04419,2.04388,2.04391,2.04522. 37 32.04819,2.06557,2.10274,2.16980,2.28060,2.4585, 38 42.742,3.20,3.96,5.30/ 39 C 40 C FBETAL IS THE LITERATURE VALUE FOR THE ECCENTRICITY 41 C CORRECTION FACTOR IN THE WALL CORRECTION EQUATION 42 C DIMENSIONLESS 43 C 44 DO 101 I=1,5 45 DO 101 J=1,3 46 READ(5,100) TEMP(I,J,1),TEMP(I,J,2),TEMP(I,J,3). 47 1TEMP(I,J,4),TEMP(I,J,5) 48 100 FORMAT(5F10.1) 49 101 CONTINUE 50 C** 

#### **FORTRAN IV (VER 45 ) SOURCE** LISTING: **88**

```
51 C TEMP IS THE MEASURED TEMPERATURE OF THE LIQUIn IN DEG.C 
 52 C 
 53 DO 103 I=1,5 
 54 DO 103 J=1,3 
           55 DO 103 K:1,5 
 56 READ(5,102) TAU(1,J,K,1),TAU(I,J,K,2),TAU(I,J,K,3). 
 57 iTAU(I,J,K,4),TAU(I,J,K,5),TAU(I,J,K,6) 
 58 102 FORMAT(6F10.1) 
         59 103 CONTINUE 
 60 C 
 61 C TAU IS THE EXPERIMENTAL SETTLING TIME/FOOT OF DISTANCE 
 62 C IN SEC/FT 
 63 C 
          64 READ(5,1C4) T8L1(1),TBL1(2),TBL1(3),T8L1(4),TBL1(5) 
 65 104 FORMAT(5I3) 
 66 DO 106 J=1.5 
 67 READ(5,105) TBL2(J,1),TBL2(J,2),TBL2(J,3),TBL2(J.4), 
 68 1T8L2(J,5) 
 69 105 FORMAT(5I3) 
 70 106 CONTINUE 
 71 READ(5,107) TBL3(1),TBL3(2),TBL3(3),TBL3(4),TBL3(5) 
          72 107 FORMAT(5I3) 
 73 C 
 74 C TBL1, TBL2 8 TBL3 ARE TABLE DESIGNATIONS IN THE OUTPUT 
 75 C 
 76 DO 1.000 [=1,5 
          77 PHIX2(I)=2.0*PHI(I) 
 78 C 
 79 C PHIX2 IS THE WEDGE ANGLE 
 80 C<br>81
     81 1000 CONTINUE 
 82 DO 1006 I=1,5 
          83 WRITE(6.1001) TBL1(I) 
 84 1001 FORMAT('1'//' ',30<sup>X</sup>,'TABLE ',12/' ',3X,'EVPERI'<br>85 1'MENTAL SETTLING TIMES FOR DELRIN SPHERES IN UC
 85 <sup>11</sup> 1'MENTAL SETTLING TIMES FOR DELRIN SPHERES IN UCON *<br>86 2'LUBRICANT'/)
 86 2'LUBRICANT'/) 
 87 WRITE(6,1002) PHIX2(I) 
 88 1002 FORMAT('-',10X,'WEDGE ANGLE= ',F5.1,' DEGREES') 
          89 DO 1006 J=1,3 
 90 WRITE(6,1003) DIAp(J) 
 91 1003 FORMAT(* */*<del>=*,1</del>0X,*SPHERE DIAMETER= *,F6.4,* INCH*<br>92 = = 1'ES*)
92 1'ES') 
          g3 WRITE(6,1004) 
 94 1004 FORMAT('-',5x,'DISTANCE FROM WEDGE',8X,'EXPERIMENTAL' 
 95 1' SETTLING TIMES'/' ',6X,'APEX TO PARTICLE',12X,'IN' 
96 2' SEC/FT OF DISTANCE'/' ',6X,'CENTER IN INCHES',6X, 
97 3'RUN 1 RUN 2 RUN 3 RUN 4 RUN 5 RUN 6'//) 
98 DO 1006 K=1,5 
99 WRITE(6,1005) X0(K),(TAU(I,J,K,L),L=1,6) 
     100 1005 FORMAT(' ',13X,F3.1,11X,6F6.1)
```

```
FORTRAN IV (VER 45 ) SOURCE LISTING: 
  101 1006 CONTINUE 
             102 DO 1009 1=1,5 
  103 DO 1009 J=1,3 
             104 DO 1009 K=1,5 
  105 SUM=0.0 
 106 DO 1007 L=1,6 
             107 UEXP(I,J,K,L)=1.0/TAU(I,J,K,L) 
  108 C 
  109 C UEXP IS THE EXPERIMENTAL SETTLING VELOCITY FOR EACH RUN 
           110 C IN A SET 
  111 C 
             112 SUM=SUM+UEXP(I,J,K,L) 
 113 1007 CONTINUE 
             114 UEXPAV(I,J,K)=SUM/6.0 
 115 C 
 116 C UEXPAV IS THE MEAN VALUE FOR THE EXPERIMENTAL VELOCITIES 
           117 C IN EACH SET 
 118 C 
 119 SUM=0.0 
 120 DO 1008 L=1'6 
             121 SUM=SUM+(UEXP(I,J,K,L)-UEXPAV(I,J,K))**2 
 122 1008 CONTINUE 
            123 SIGMA(I,J,K)=SQRT(SUM/6.0) 
 124 C 
 125 C SIGMA IS THE STANDARD DEVIATION FOR EACH SET 
 126 C 
 127 1009 CONTINUE 
 128 DO 1017 I=1,5<br>129 DO 1017 K=1,5
 129 DO 1017 K=1,5 
 130 WRITE(6.1010) TBL2(I,K) 
      131 1010 FORMAT('1'//' ',30X,'TABLE ',I2/'1X,'EXPERIMENT' 
 132 1/AL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON' 
 133 2' LUBRICANT'/) 
            134 WRITE(6,1011) PHIX2(I),X0(K) 
 135 1011 FORMAT('0',10X,'WEDGE ANGLE= ',F5,1,' DEGREES'/' '<br>136 110X,'DISTANCE FROM WEDGE APEX TO PARTICLE CENTER=
           136 110X,'DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= ', 
 137 2F3.1,' INCHES') 
 138 DO 1017 J=1,3 
            /39 WRITE(6,1012) DIAP(J) 
 140 1012 FORMAT('—',,l0X,'SPHERE DIAMETER: ',F6,4,' INCHES') 
 141 WRITE(6.1013) 
 142 1013 FORMAT('—',9X,'RUN NUMBER',19X,'EXPERIMENTAL PARTIC' 
 143 1'LE'/' ',11X,'IN SET',17X,'SETTLING VELOCITIES TN 
           144 2'FT/SEC'/) 
 145 DO 1015 L=1,6 
 146 WRITE(6,1014) L,UEXP(I,J,K,L)<br>147 1014 FORMAT(* *,13X,I1,29X,F11.9)
       147 1014 FORMAT(' ',13x,I1,29x,F11,9) 
 148 1015 CONTINUE 
            149 WRITE(6,1016) UEXPAV(I,J,K),SIGMA(I,J,K) 
 150 1016 FORMAT(' ',4X,'MEAN VALUE OF SET IS',19X,F11.9/''
```
 $\rightarrow$ 

**FORTRAN IV (VER 45 ) SOURCE** LISTING: **90 151 12X,'WITH STANDARD DEVIATION** OF',15x,F11.9 **152. 1017 CONTINUE 153 DO 1020 I=1,5 154 DO 1020 J=1.3 155 DO 1020 K=1.5 156 IF(FPHI1(I).EQ.TEST.OR.FPHI2(1).EQ.TEST) GO TO 1020 157 RHOL(I,J,K)=1.076667-0.75889E-03\*TEMP(I,J,K) 158 C 159 C** RHOL **IS THE CALCULATED DENSITY OF THE LIQUID IN GMS/CC 160 C 161 MUKL(I,J,K)=EXP(EXP(5.138245-0.561074\*(ALOG(1.8\*TEMP 162 1(I,J,K)+132.0)))-1.7) 163 C 164 C MUKL IS THE CALCULATED KINEMATIC VISCOSITY OF THE LIQUID 165 C IN CENTISTOKES 166 C 167 MUAL(I,J,K)=MUKL(I,J,K)\*1.076391E-05\*RHOL(I,J,K)\* 168 162.42642 169 C 170 C MUAL IS THE CALCULATED ABSOLUTE VISCOSITY OF THE LIQUID 171 C IN LB/FT-SEC 172 C 173 US(I,J,K)=G\*((DIAP(J)/12.0)\*\*2)\*(RHOP(J)-RHOL(I,J,K) 174 1)\*62.42642/(18.0\*MUAL(I,J,K)) 175 C 176 C US IS THE STOKES SETTLING VELOCITY OF THE PARTICLE IN 177 C FT/SEC 178 C 179 U1(I,J,K)=US(I,J,K)\*(1.0-DIAP(J)/2.0\*FPHI1(I)/xo(K)- 180 1(((DIAP(J)/2.0)/xo(K))\*\*3)\*FPHI2(I)) 181 C 182 C U1 IS THE CALCULATED SETTLING VELOCITY OF THE** PARTICLE **183 C USING THE WEDGE CORRECTION IN FT/SEC 184 C 185 DO 1018 11=2,26 186 IF(X0(K)/RO,LE.BETAL(II).AND.XO(K)/RO.GE.BETAL(II-1 187 1)) GO TO 1019 188 1)18 CONTINUE 189 1019 BETA1=BETAL(II-1) 190 BETA2=BETAL(II) 191 FBETA1=FBETAL(II-1) 192 FBETA2=FBETAL(II) 193 FBETA(I,J,K)=(((XO(K)/RO-BETA1)/(BETA2-BETA1))\*( 194 1FBETA2-FBETA1))+FBETA1 195 C 196 C FBETA IS THE INTERPOLATED ECCENTRICITY CORRECTION FACTOR 197 C IN THE WALL CORRECTION EQUATION DIMENSIONLESS 198 C 199 U2(I,J,K)=US(I,J,K)\*(1.0-DIAP(J)/2.0\*FPHI1(I)/XO(K)- 200 1((DIAP(J)/2.0)/XO(K))\*\*3)\*FPHI2(I)-DIAP(J)/(2.0\*RO)** 

#### **FORTRAN IV (VER 45 ) SOURCE LISTING:**

```
201 2*(PHI(I)/180.0)*FBETA(I,J,K)) 
202 C 
203 C U2 IS THE CALCULATED SETTLING VELOCITY OF THE PARTICLE 
         204 C USING THE WEDGE AND MODIFIED WALL CORRECTION IN FT/SEC 
205 C 
206 1023 CONTINUE 
207 DO 1028 I=1.5 
208 WRITE(6,1021) TBL3(I) 
209 1021 FORMAT('1'//',30X,'TABLE ''I2/'',2x.'CALCULATED' 
210 1' SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON ' 
211 2'LUBRICANT'//)<br>212 RRITE(6,1022)
212 WRITE(6,1022) PHIX2(I) 
     213 1022 FORMAT('u',10X,'WEDGE ANGLE= ',F5.1,' DEGREES') 
214 DO 1028 J=1,3 
215 WRITE(6,1023) DIAP(J) 
216 1023 FORMAT(' '/'-',10X,'SPHERE DIAMETER= ',F6.4,' INCH<sup>'</sup><br>217 1'ES')
217 1'ES') 
218 WRITE(6,1024) 
219 1024 FORMAT('—',3X,'DISTANCE FROM WEDGE',7X,'CALCULATED' 
220 1' PARTICLE SETTLING VELOCITIES'/' ',4X,'APEX TO ' 
         221 2'PARTICLE',9X,'IN FT/SEC. CORRECTED FOR THE EFFECTS' 
222 3' OF'/' ',4X,'CENTER IN INCHES',14X,'WEDGE ONLY',2X. 
223 4'||||',2X,'WEDGE 8 WALL'//)<br>224 . DO 1028 K=1,5
224 00 1028 K=1.5 
225 IF(FPHI1(I),EQ.TEST.OR'.FPHI2(I).EQ.TEST) GO TO 1026 
226 WRITE(6.1025) X0(K),U1(I,J,K),U2(I,J,K) 
227 1025 FORMAT(' ',11X,F3.1.19X,F11.9,9X,F11.9) 
228 GO TO 1028 
     229 1026 WRITE(6,1027) XO(K) 
230 1c27 FORMAT(' ',11X, F3.1,18X, 'NOT AVAILABLE', 7X, 'NOT AVA'
231 1'ILABLE') 
     232 1028 CONTINUE 
233 WRITE(6.1029) 
     234 102i) FORMAT('1') 
235 STOP 
236 END
```
# APPENDIX E

 $\sim 100$ 

 $\sim 10^7$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$ 

## **TAM 11**

EXPERIMENTAL SETTLING **TIMES FOR DELRIN SPHERES IN UCON LUBRICANT** 

 $\sim$   $\zeta$  .

 $\sim 100$ 

**WEDGE** ANGLE:: 60.0 **DEGREES** 

 $\sim$ 

SPHERE DIAMETER= 0.1562 INCHES



#### **SPHERE DIAMETER= 0.2497 INCHES**



#### **SPHERE DIAMETER= 0.3435 INCHES**



 $\sim$ 

 $\sim$   $\alpha$ 

# TABLE 12

EXPERIMENTAL SETTLING TIMES FOR DELRIN SPHERES IN UCON LUBRICANT

WEDGE ANGLE= 90.0 DEGREES

## SPHERE DIAMETER= 0.1562 INCHES



#### SPHERE DIAMETER= 0.2497 INCHES



## SPHERE DIAMETER= 0.3435 INCHES



 $\mathbf{f}$ 

 $\sim 10$ 

 $\hat{\vec{r}}$ 

 $\mathcal{L}$  $\mathcal{L}_{\rm{max}}$
### TABLE 13 EXPERTAL SETTLING TIMES FOR DELRIN SPHERES IN UCON LUBRICANT

### WEDGE ANGLE= 180.0 DEGREES

### SPHERE DIAMETER= 0.1562 INCHES



#### SPHERE DIAMETER= 0,2497 INCHES

 $\sim$ 





 $\sim$ 

#### SPHERE DIAMETER= 0.3435 INCHES



 $\bullet$ 

 $\mathcal{A}^{\mathcal{A}}$ 

EXPERIMENTAL SETTLING TIMES FOR DELRIN SPHERES IN UCON LUBRICANT

WEDGE ANGLE= 270.0 DEGREES

### SPHERE DIAMETER= 0.1562 INCHES



#### SPHERE DIAMETER= 0.2497 INCHES

 $\sim 10$ 



### SPHERE DIAMETER= 0,3435 INCHES



 $\sim$   $\sim$ 

WEDGE ANGLE= 360.0 DEGREES

### SPHERE DIAMETER= 0.1562 INCHES



 $\sim$ 

### SPHERE DIAMETER= 0.2497 INCHES

 $\mathbb{D}$ 



### SPHERE DIAMETER= 0.3435 INCHES



APPENDIX F

 $\bar{z}$ 

 $\bar{z}$ 

 $\sim$   $\sim$ 

 $\hat{\mathcal{L}}$ 

 $\sim$   $\sim$ 

EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 60.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 1.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER EXPERIMENTAL PARTICLE<br>IN SETTE SETTLING VELOCITIES IN FT. SETTLING VELOCITIES IN FT/SEC  $\begin{array}{c} 1 \\ 2 \\ 3 \end{array}$   $\begin{array}{c} 0.002762431 \\ 0.002751789 \\ 0.002743484 \end{array}$ 2 0.002751789  $\frac{3}{4}$  0.002743484<br>4 0.002749519 4 0.002749519<br>5 0.002757859 5 0.002757859 0.002750274<br>0.002752559 MEAN VALUE OF SET IS 0.002752559<br>TH STANDARD DEVIATION OF 0.00006096

TABLE 16

SPHERE DIAMETER= 0.2497 INCHES

WITH STANDARD DEVIATION OF

 $\bar{\mathbf{v}}$ 



SPHERE DIAMETER= 0.3435 INCHES



 $\sim$ 

TABLE 17 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

> WEDGE ANGLE= 60.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 2.0 INCHES

SPHERE DIAMETER= 0.1562 INCHES

 $\ddot{\phantom{a}}$ 

 $\Delta$ 

RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC



SPHERE DIAMETER: 0.2497 INCHES



SPHERE DIAMETER: 0.3435 INCHES



TABLE 18 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

WEDGE ANGLE= 60.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 3.0 INCHES

SPHERE DIAMETER= 0.1562 INCHES

 $\epsilon$ 



SPHERE DIAMETER= 0.2497 INCHES



SPHERE DIAMETER= 0,3435 INCHES

 $\bar{a}$ 



TARLE 19 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 60.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 4.0 INCHES SPHERE DIAMETER= 0.1562 INCHES EXPERIMENTAL PARTICLE RU' NUMBER SETTLING VELOCITIES IN FT/SEC IN SET 1 0.003086420<br>2 0.003080714 2<br>2<br>3 0.003087373 3 0.003087373 4 0.003091190<br>5 (a) 0.003084516 5 0.003084516 6 0.003088325 MEAN VALUE OF SET IS 0.003086422 WITH STANDARD DEVIATION OF  $\mathcal{A}$ SPHERE DIAMETER= 0.2497 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET  $\frac{1}{2}$  0.007423904<br>0.007462684  $2$ <br>3 0.007462684<br>0.007473841  $\begin{array}{ccc} 3 & 0.007473841 \ 4 & 0.007412896 \end{array}$ 4 0.007412896<br>5 0.007440474  $\begin{array}{c|c} 5 & 0.007440474 \\ 6 & 0.007429417 \end{array}$ 0.007429417<br>0.007440533 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF TH 0.000021463 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.3435 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET  $\frac{1}{2}$  0.014619883<br>0.014814813  $2$ <br>3 0.014814813<br>0.014727540  $\frac{3}{4}$  0.01472754 $\frac{1}{0}$ 4 0.014771048<br>5 0.014684286  $\begin{array}{ccc} 5 & 0.014684236 \\ 6 & 0.014705881 \end{array}$ 0.014705881<br>0.014720559 MEAN VALUE OF SET IS 0.014720559 WITH STANDARD DEVIATION OF

 $\mathcal{A}^{\pm}$ 

# TABLE 20 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 60.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 5.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.003112355<br>0.003104625  $\frac{2}{3}$  0.003104625<br>3 0.003109452  $\frac{3}{4}$  0.003109452 4 0.003111389<br>5 0.003109452 5 0,003109452 0.003107520<br>0.003109131 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF 0.000002536 WITH STANDARD DEVIATION OF

SPHERE DIAMETER= 0,2497 INCHES



SPHERE DIAMETER= 0.3435 INCHES



**TABLE** 21 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

> WEDGE ANGLE= 90.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 1.0 INCHES

SPHERE DIAMETER= 0.1562 INCHES



SPHERE DIAMETER= 0.2497 INCHES



SPHERE DIAMETER= 0.3435 INCHES

 $\sim 10^{-11}$ 



 $\sim$ 

 $\sim 20$ 

 $\sim 10^{-10}$ 

TABLE 22 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 90.0 DEGREES DISTANCE FROM 'WEDGE APEX TO PARTICLE CENTER= 2.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.003052502<br>0.003046923  $\frac{2}{3}$ <br>3  $\frac{3}{4}$  0.003053435<br>4 0.003058104 0.003058104 5 0.003052502 6 0.003051572<br>F OF SET IS 0.003052505 MEAN VALUE OF SET IS 0.003052505 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.007347535  $\frac{2}{3}$  0.007336754<br>3  $\frac{3}{4}$  0.007299267<br>4 0.007283319  $\frac{1}{4}$  0.007283319<br>5 0.007304601 5 0.007304601<br>6 0.007278018 0.007278018<br>0.007308248 MEAN VALUE OF SET IS<br>TH STAMDARD DEVIATION OF 0.000025773 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.3435 INCHES RUN NUMBER **EXPERIMENTAL PARTICLE** IN SET SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.0143471.99  $\frac{2}{3}$  0.014285713  $\begin{array}{ccc} 3 & 0.014285713 \\ 4 & 0.014367811 \end{array}$ 4 0.014367811<br>5 0.014224749 5 0.014224749 ة 0.014245015<br>For SET IS = 0.014292687 MEAN VALUE OF SET IS **0.014292687**<br>TH STANDARD DEVIATION OF **0.000050991** 

WITH STANDARD DEVIATION OF

**TABLE 23 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 90.3 CEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER: 3.0 INCHES SPHERE DIAMETER= 0:1562 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET**  1 **0.003100775**   $\frac{2}{3}$  0.003105590<br>3 3 0.003095016<br>4 0.003098854 4 0.003098854<br>5 0.003097893 5 0.003097893<br>6 0.003101738 0.003101738<br>0.003099977 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF **1888 0.000003303** WITH STANDARD DEVIATION OF **SPHERE DIAMETER= 0.2497 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN** FT/SEC IN SET **1 0.007513147 2** 0.007501874 3 0.007457118<br>4 0.007490635 **<sup>4</sup>**0.007490635 0.007507507 6<br>
E OF SET IS<br>
E OF SET IS<br>
O.007500064 MEAN VALUE **OF SET IS 0.007500064**  WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.3435 INCHES **RUN NUMBER EXPERIMENTAL PARTICLE**  IN SET **SETTLING VELOCITIES IN FT/SEC**  $\frac{1}{2}$  0.015015014<br>0.014970057 2 0.014970057<br>3 0.014836796 **3 0.014836796 4** 0.014749259 5 **0.014749259 6 0.014836796 MEAN VALUE OF SET IS** 0.014839527 WITH STANDARD DEVIATION OF

 $\sim 100$ 

TABLE 24 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 90.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 4.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.003121099<br>0.003124023  $\frac{2}{3}$  0.003124023<br>3  $\frac{3}{4}$  0.003118179<br>4 0.003126954 **4** 0.003126954 5 0.003127933<br>6 0.003121099 0.003121099<br>0.003123213 MEAN VALUE OF SET IS 0.003123213<br>TH STANDARD DEVIATION OF 0.000003445 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES RUM NUMBER IN SET EXPERIMENTAL PARTICLS SETTLING VELOCITIES IN FT/SEC 1  $\frac{2}{3}$ 0.007598784 0.007616144  $\frac{3}{4}$  0.007598784<br>4 0.007570021  $\frac{4}{5}$  0.007570021 5 0.007570021 0.007604562<br>0.007593051 MEAN VALUE OF SET IS 0.007593051 WITH STAMDARD DEVIATION OF SPHERE DIAMETER= 0,3435 INCHES RUN NUMBER **EXPERIMENTAL PARTICLE**<br>IN SET **ETTLING VELOCITIES IN FT** SETTLING VELOCITIES IN FT/SEC 1<br>2 0.015243899<br>0.015128590  $\frac{2}{3}$  0.01512859 $\frac{1}{0}$  0.01512859 $\frac{1}{3}$ 3<br>4 0.015151512<br>9.015151512 4 0.015151512<br>5 0.015128590 5 0.015128590 0.015151512<br>0.015140273 MEAN VALUE OF SET IS 0.015140273 WITH STANDARD DEVIATION OF

107

 $\sim$   $\lambda$ 

TABLE 25 EXPERITAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 90.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER: 5.0 INCHES SPHERE DIAMETER: 0.1562 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET 1 0.00314169ñ<br>2 0.003137748 2<br>2<br>3 0.003134796  $\begin{array}{ccc} 3 & 0.003134796 \\ 4 & 0.003138731 \end{array}$ 4 0.003138731<br>5 0.003139717 5 0.003139717<br>6 0.003139717 0.003139717<br>0.003138731 MEAN VALUE OF SET IS 0.003138731 WITH STANDARD DEVIATION OF SPHERE DIAMETER: 0.2497 INCHES RUN NUMBER EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC IN SET 0.00768639i 1 2 0.007633585<br>3 0.007680491 3 0.007680491<br>4 0.007662833 4 0.007662833<br>5 0.007662833 5 0.007662833 0.007656965<br>0.007663850 MEAN VALUE OF SET IS 0.007663850 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.3435 INCHES



### **TABLE 26 EXPERIMENTAL SETTLING VELOCITIES FOR** DELRIN SPHERES TN UC1N LUBRICANT

WEDGE ANGLE= **180.0 DEGREES DISTANCE FROM WEDGE Apex TO PARTICLE CENTER: 1.0 INCHES** 

**SPHERE DIAMETER= 0,1562 INCHES** 



**SPHERE DIAMETER= 0.2497 INCHES** 



SPHERE DIAMETER= 0.3435 INCHES



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TABLE 27 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

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WEDGE ANGLE= 180.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 2.0 INCHES

SPHERE DIAMETER= 0.1562 INCHES

RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.003121099<br>0.003125978 2<br>2<br>3 0.003125978  $\frac{3}{4}$  0.003125978<br>4 0.003120125 4 0.003120125<br>5 0.003121099 5 0.003121099<br>6 0.003118179 0.003118179<br>0.003122075 MEAN VALUE OF SET IS 0.003122075<br>TH STANDARD DEVIATION OF 0.000002925 WITH STANDARD DEVIATION OF

DIAMETER= 0.2497 INCHES



SPHERE DIAMETER= 0.3435 INCHES

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TABLE 28 EXPERIMENTAL SETTLING VELOCITIES FOR CELRIN SPHERES TN UCON LUBRICANT

> WEDGE ANGLE= 180.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 3.0 INCHES

SPHERE DIAMETER= 0.1562 INCHES



SPHERE DIAMETER= 0.2497 INCHES



SPHERE DIAMETER= 0,3435 INCHES



TABLE 29 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCOM LUBRICANT WEDGE ANGLE= 180.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 4.0 INCHES

SPHERE DIAMETER= 0.1562 INCHES



SPHERE DIAMETER= 0.2497 INCHES



SPHERE DIAMETER= 0,3435 INCHES



EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 180.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 5.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET **EXPERIMENTAL PARTICLE** SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.00316556 $\frac{1}{0}$  0.003165556  $\frac{2}{3}$  0.003163555<br>0.003166561  $\frac{3}{4}$  0.003166561<br>0.003168567 4 0.003168567<br>5 0.003165560 0.003165560 6 0.003165566<br>FOR SET IS 0.003165892 MEAN VALUE OF SET IS 0.003165892 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES

TABLE 30



#### SPHERE DIAMETER= 0.3435 INCHES

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RUN NUMBER EXPERIMENTAL PARTICLE<br>IN SET SETTLING VELOCITIES IN FT. SETTLING VELOCITIES IN FT/SEC 1 2 3 4 5 6 MEAN VALUE OF SET IS WITH STANDARD DEVIATION OF 0.015673986 0.015625000 0.015600622 0.015625000 0.015600622 0.015625000 0.015625030 0.000024453

 $\mathcal{A} \subset \mathcal{F}$  .

TABLE 31 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 270.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 1.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC 1 0.003070310 2<br>2<br>3 0.003067485  $\frac{3}{4}$ <br>4 0.003068427 4 0.003068427<br>5 0.003076923 5 0.003076923<br>6 0.003065604 0.003065604<br>0.003068275 MEAN VALUE OF SET IS 0.003068275 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.007385522<br>0.007390980  $\frac{2}{3}$ <br>3 0.007390980<br>0.007352941  $\begin{array}{ccc} 3 & 0.007352941 \\ 4 & 0.007374629 \end{array}$ 4<br>5<br>5<br>0.007374629 5 0.007374629 0.007407404<br>0.007381015 MEAN VALUE OF SET IS WITH STANDARD DEVIATION O SPHERE DIAMETER= 0.3435 INCHES **0.000016762**  RUN NUMBER **IN SET**  EXPERIMENTAL PARTICLE **SETTLING VELOCITIES IN FT/SEC 1 0.014641285 2 0.014598537 3 0.014598537 4 0.014492756**  0.014684286 6 **0.014662754** 

**MEAN VALUE** OF SET IS 0.014613021

WITH STANDARD DEVIATION OF

TABLE 32 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 270.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 2.0 INCHES SPHERE DIAMETER= 0.1562 INCHES EXPERIMENTAL PARTICLE RUN NUMBER IN SET SETTLING VELOCITIES IN FT/SEC 0.003123048 1  $\frac{2}{3}$  0.003131850<br>0.003127933  $\frac{3}{4}$  0.003127933<br>0.003134796 4 0.003134796<br>5 0.003127933 5 0.003127933 6 0.003126954<br>F OF SET IS 0.003128751 MEAN VALUE OF SET IS 0.003128751 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET  $\frac{1}{2}$  0.007598784<br>0.007633585  $\frac{2}{3}$  0.007633585<br>5 0.007651109  $\frac{3}{4}$  0.007651109 4 0.007621948<br>5 0.007621948 5 0.007621948<br>6 0.007627763 0.007627763<br>0.007625856 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF 0.000015623 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.3435 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET  $\frac{1}{2}$  0.015243899<br>0.015174508  $\frac{2}{3}$  0.015174508<br>3 0.015290521  $\begin{array}{ccc} 3 & 0.015290521 \\ 4 & 0.015220698 \end{array}$ 4 0.015220698<br>5 0.015313935 0.015313935 6 0.015360981<br>F OF SET IS 0.015267409 MEAN VALUE OF SET IS<br>TH STARDARD DEVIATION OF **1888 0.000061671** WITH STANDARD DEVIATION OF

TABLE 33 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 270.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 3.0 INCHES

EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET  $\frac{1}{2}$  0.003142678<br>0.003150600  $\frac{2}{3}$  0.003150600<br>0.003145644  $\frac{3}{4}$  0.003145644<br>4 0.003151591  $\frac{4}{5}$ <br>5 5 0.003153578<br>6 0.003148613 0.003148613<br>0.003148782 MEAN VALUE OF SET IS 0.003148782 WITH STANDARD DEVIATION OF

SPHERE DIAMETER= 0.2497 INCHES

SPHERE DIAMETER= 0.1562 INCHES

RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.007704157<br>0.007727973  $\frac{2}{3}$  0.007727973<br>0.007716049  $\frac{3}{4}$  0.007716049 4 0.007704157<br>5 0.007692307 5 0.007692307 0.007674593<br>0.007703204 MEAN VALUE OF SET IS 0 0.007703204<br>TH STAMDARD DEVIATION OF 0.000016919 WITH STANDARD DEVIATION OF

SPHERE DIAMETER= 0.3435 INCHES

RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC 1. 0.015384614<br>2. 0.015479874 2 0.015479874<br>3 0.015503876  $\frac{3}{4}$  0.015503876<br>4 4 0.015552096<br>5 0.015479874 **5** 0.015479874 6 0.015479874 MEAN VALUE OF SET IS 0.015480030 WITH STANDARD DEVIATION OF

 $\sim$ TABLE 34 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

> WEDGE ANGLE= 270.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 4.0 INCHES

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SPHERE DIAMETER= 0.1562 INCHES



SPHERE DIAMETER= 0.2497 INCHES



SPHERE DIAMETER= 0.3435 INCHES



TABLE 35 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

> WEDGE ANGLE= 270.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER: 5.0 INCHES

SPHERE DIAMETER= 0.1562 INCHES



SPHERE DIAMETER= 0.2497 INCHES

RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC 1 0.00775193A 0.007739935  $\frac{3}{4}$  0.007757951<br>4 0.007745933 4 0.007745933<br>5 0.007776048 5 0.007776048<br>6 0.007757951 0.007757951<br>0.007754959 MEAN VALUE OF SET IS CONTROL 10.007754959<br>TH STANDARD DEVIATION OF COORD 0.000011394 WITH STANDARD DEVIATION OF

SPHERE DIAMETER= 0.3435 INCHES



TABLE 36 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES. IN UCON LUBRICANT WEDGE ANGLE= 360.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 1.0 INCHES **SPHERE DIAMETER= 0.1562 INCHES**  RUN NUMBER IN SET **EXPERIMENTAL PARTICLE SETTLING VELOCITIES** IN FT/SEC **1 0.002988642 2 0.002994012 3 0.002985967 4 0.002979737 5 0.002987751 6 0.002984183 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF <b>1888** 0.000004353 WITH STANDARD DEVIATION OF **SPHERE DIAMETER= 0:2497** INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.007042252  $\frac{2}{3}$ <br>3 0.007062145 **3 0.007062145 4** 0.007087171  $\frac{5}{6}$  0.00710227 $\frac{3}{10}$ 0.007087171<br>0.007071368 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF **1888 0.000022251** WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.3435 INCHES RUN NUMBER <sup>1</sup>'4 SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.014124293<br>0.01388888  $\frac{2}{3}$  0.013888888<br>0.014084507  $\frac{3}{4}$  0.014084507<br>4 0.014005601 4 0.014005601<br>5 0.014064696 0.014064696 6 0.014044944<br>F OF SET IS 0.014035482 MEAN VALUE OF SET IS 0.014035492

WITH STANDARD DEVIATION OF

TABLE 37 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 360.0 DEGREES DISTANCE FRON WEDGE APEX TO PARTICLE CENTER= 2.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.003041362<br>0.003039513 2 0.003039513 3  $\frac{4}{5}$ 0.003032140 0.003033061 0.003033980 6<br>
F OF SET IS<br>
F OF SET IS 0.003036131 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF **1888 0.000003395** WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES RUN NUMBER **EXPERIMENTAL PARTICLE**<br>IN SET SETTLING VELOCITIES IN FT SETTLING VELOCITIES IN FT/SEC  $\frac{1}{2}$  0.007251631<br>0.007246375 2 0.007246375<br>3 0.007267438 3 0.007267438 4 0.007267438<br>5 0.007251631 5 0.007251631 0.007299267<br>0.007263962 EAN VALUE OF SET IS 0.007263962 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.3435 INCHES RUN NUMBER EXPERIMENTAL PARTICLE<br>IN SET SETTLING VELOCITIES IN FT. SETTLING VELOCITIES IN FT/SEC

WITH STANDARD DEVIATION OF

 $\frac{1}{2}$  0.014367811<br>0.014513787  $2$ <br>3 0.014513787<br>0.01449275ñ  $\frac{3}{4}$  0.014492750 4 0.014534879<br>5 0.014471777 5 0.014471777<br>6 0.014471777 0.014471777<br>0.014475454 MEAN VALUE OF SET IS 0.014475454

TABLE 38 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 360.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 3.0 INCHES SPHERE DIAMETER= 0.1562 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET 1 0.003058104 0.00304878n 2 0.003054369 3 4 0.003058104<br>5 0.003055300 5 0.003055300<br>6 0.003057168 0.003057168<br>0.003055302 MEAN VALUE OF SET IS 0.003055302 WITH STANDARD DEVIATION OF SPHERE DIAMETER: 0.2497 INCHES EXPERIMENTAL PARTICLE RUN NUMBER SETTLING VELOCITIES IN FT/SEC IN SET 0.007347535 1 0.007331375  $\frac{2}{3}$ 3<br>4 0.007304601  $\frac{4}{5}$  0.007304601<br>5  $\begin{array}{ccc} 5 & 0.007331375 \\ 6 & 0.007342141 \end{array}$ 0.007342141<br>0.007333193 MEAN VALUE OF SET IS<br>TH STAMDARD DEVIATION OF **1988** 0.000014083 WITH STANDARD DEVIATION OF

SPHERE DIAMETER= 0.3435 INCHES



**TABLE 39 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN** UCON LUBRICANT **WEDGE ANGLE= 360.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER: 4.0 INCHES**  SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC 1<br>2<br>2<br>0.003063726  $\frac{2}{3}$  0.003063726<br>0.003065604 3<br>4 0.003067485<br>4 0.003067485 4 0.003067485<br>5 0.003062787 5 0.003062787<br>6 0.003063726 0.003063726<br>0.003063882 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF **18900002326** WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES **RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC**   $\frac{1}{2}$  0.007352941<br>0.007358350 **2** 0.00735835n **3 0.007352941 4 0.007352941**  5 0.007336754<br>6 0.007331375 **6 0.007331375 MEAN VALUE** OF SET **IS 0.007347550**  WITH STANDARD DEVIATION OF SPHERE **DIAMETER=** 0:3435 INCHES RUN NUMBER **SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN** FT/SEC 1 0.014684286<br>2 0.014641285 2 0.014641285<br>3 0.014705881  $\frac{3}{4}$  0.014705881<br>4 0.014792897 4 0.014792897<br>5 0.014727548

MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF **10.000046808** WITH STANDARD DEVIATION OF

**5 0.014727540**  0.014684286<br>0.014706016 **122** 

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TABLE 40 EXPERIMENTAL SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT WEDGE ANGLE= 360.0 DEGREES DISTANCE FROM WEDGE APEX TO PARTICLE CENTER= 5.0 INCHES SPHERE DIAMETER= 0.1562 INCHES RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC 1  $\frac{2}{3}$ 0.003066543 0.003065604  $\frac{3}{4}$  0.003068427<br> $\frac{3}{4}$ 0.003067485 5 0.003069367<br>6 0.003070310 0.003070310<br>0.003067954 MEAN VALUE OF SET IS 0.003067954 WITH STANDARD DEVIATION OF SPHERE DIAMETER= 0.2497 INCHES RUN! NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\begin{array}{c} 1 \\ 2 \end{array}$   $\begin{array}{c} 0.007352941 \\ 0.007352941 \end{array}$  $\frac{2}{3}$ <br>3<br>3  $\frac{3}{4}$ <br>4 0.007380072  $\frac{1}{4}$ <br>5 0.007374629 5 0.007374629 0.007380072<br>0.007366501 MEAN VALUE OF SET IS<br>TH STANDARD DEVIATION OF 1000012032

SPHERE DIAMETER= 0.3435 INCHES

WITH STANDARD DEVIATION OF

RUN NUMBER IN SET EXPERIMENTAL PARTICLE SETTLING VELOCITIES IN FT/SEC  $\begin{array}{ccc} 1 & 0.014684286 \\ 2 & 0.014749259 \end{array}$  $\frac{2}{3}$ <br>3 0.014771048  $\frac{3}{4}$  0.014771048 4 0.014792897<br>5 0.014727540  $\frac{5}{6}$  0.014727546<br>6 0.014749259 0.014749259<br>0.014745701 MEAN VALUE OF SET IS 0.014745701 WITH STANDARD DEVIATION OF

APPENDIX G

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TABLE 41 CALCULATED SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

WEDGE ANGLE= 60.0 DEGREES

 $\sim$   $\sim$ 

 $\mathcal{A}^{\pm}$ 

SPHERE DIAMETER= 0.1562 INCHES





#### SPHERE DIAMETER= 0,2497 INCHES





### SPHERE DIAMETER= 0.3435 INCHES

 $\mathcal{L}(\mathcal{L})$  and  $\mathcal{L}(\mathcal{L})$  .



 $\sim 0.5$ 

 $\sim 10^{11}$  m  $^{-1}$ 

 $\sim 10^{-11}$ 

TABLE 42 CALCULATED SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

 $\rightarrow$ 

WEDGE ANGLE= 90.0 DEGREES

SPHERE DIAMETER= 0.1562 INCHES

DISTANCE FROM WEDGE CALCULATED PARTICLE SETTLING VELOCITIES<br>APEX TO PARTICLE TN FT/SEC. CORRECTED FOR THE EFFECTS OF APEX TO PARTICLE IN FT/SEC. CORRECTED FOR THE EFFECTS OF CENTER IN INCHES WEDGE ONLY 1111 WEDGE & WALL



#### SPHERE DIAMETER= 0.2497 INCHES

 $\sim 10^{-10}$ 

 DISTANCE FROM WEDGE APEX TO PARTICLE CENTER IN INCHES CALCULATED PARTICLE SETTLING VELOCITIES IN FT/SEC. CORRECTED FOR THE EFFECTS OF WEDGE ONLY ITTI WEDGE-8 WALL



SPHERE DIAMETER= 0.3435 INCHES



 $\sim 10^6$ 

**Contract** 

TABLE 43

CALCULATED SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

WEDGE ANGLE= 180.0 DEGREES

SPHERE DIAMETER= 0.1562 INCHES



 $\mathcal{L}^{\text{max}}_{\text{max}}$  ,  $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\sim 10^7$ 



#### SPHERE DIAMETER= 0.2497 INCHES





### SpHERE DIAMETER= 0.3435 INCHES



TABLE 44

 $\sim 100$ 

CALCULATED SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

WEDGE ANGLE= 270.0 DEGREES

 $\sim$ 

#### SPHERE DIAMETER= 0.1562 INCHES

DISTANCE FROM WEDGE CALCULATED PARTICLE SETTLING VELOCITIES<br>APEX TO PARTICLE TN FT/SEC. CORRECTED FOR THE EFFECTS OF APEX TO PARTICLE IN FT/SEC. CORRECTED FOR THE EFFECTS OF CENTER IN INCHES WEDGE ONLY ITTT WEDGE & WALL



#### SPHERE DIAMETER= 0.2497 INCHES



NOT AVAILABLE

#### SPHERE DIAMETER= 0.3433 INCHES



TABLE 45

CALCULATED SETTLING VELOCITIES FOR DELRIN SPHERES IN UCON LUBRICANT

WEDGE ANGLE= 360.0 DEGREES

 $\Delta \Delta$ 

 $\lambda$ 

SPHERE DIAMETER: 0.1562 INCHES





4.0 0.003175199<br>5.0 0.003181175





SPHERE DIAMETER= 0.3435 INCHES



0.003138531

## NOMENCLATURE



 $\bar{\alpha}$ 

 $\bar{\epsilon}$ 

 $\mathbb{Z}^2$
- = eccentricity ratio  $(x_0/R_0)$ , dimensionless  $\beta$ = viscosity, lb/ft-sec  $\mathfrak{u}$  $\rho_1, \rho_p$  = density, lb/ft<sup>3</sup> a = standard deviation for experimental settling velocities  $\tau$  = experimental particle settling time, sec  $\phi$  = half of the wedge angle, degrees
- w = angular velocity, rad/sec

Subscripts

- e = experimental
- em = experimental mean
- $1 = liquid$
- $o = center$
- p = particle
- s = Stokes
- 1,2 = reference subscripts for constants

## LITERATURE REFERENCES

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